

CALIBRATION OF A 200 FT.  
WAVE TANK AND A STUDY  
OF THE RESPONSE OF A  
CIRCULAR CYLINDER TO  
SINUSOIDAL WAVES

CENTRE FOR NEWFOUNDLAND STUDIES

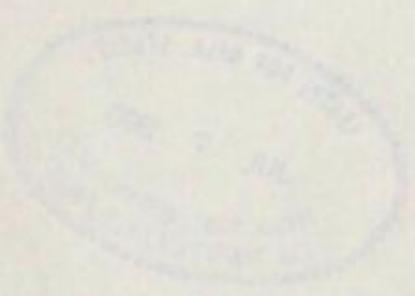
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CALIBRATION OF A 200 ft WAVE TANK AND A STUDY  
OF THE RESPONSE OF A CIRCULAR CYLINDER TO SINUSOIDAL WAVES

by

Chun-Ming Chen, B.Sc.



A Thesis submitted in partial fulfillment of the  
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### ABSTRACT

A series of experiments have been carried out to verify linear wave generator theory as well as to study wave forces on a vertical circular cylinder. The experimental data shows that linear wave generator theory can accurately predict the relationship between water depth, stroke of the wave generator, wave height, and wave period. Experimental forces and overturning moments on the cylinder show reasonably good agreement with linear wave theory and Morison's equation.

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### NOMENCLATURE

a	amplitude of the incident wave
$a_1$	modal coefficient
A	projected area perpendicular to the stream velocity
b	width of the wave tank
$C_D$	drag coefficient
$C_L$	lift coefficient
$C_M$	mass coefficient
d	still water depth
$d/L$	relative depth
D	cylinder diameter
$D/L$	relative diameter
e	stroke of the wave generator (one-half its total motion)
E	Elastic modulus
f	wave frequency, $1/T$
$f_1$	fundamental frequency of the cylinder
$f_i$	natural frequencies of the cylinder
$f_v$	vortex shedding frequency
$F_D$	drag force
$F_I$	inertia force
$F_L$	lift force
$F_h(s)$	horizontal longitudinal force on a differential length ds of the cylinder
$F_h$	total horizontal longitudinal force

$F_T(s)$	horizontal transverse force on a differential length ds of the cylinder
$F_T$	total horizontal transverse force
$g$	acceleration due to gravity
$H$	wave height
$H/L$	wave steepness
$I$	moment of inertia
$k$	wave number, $2\pi/L$
$K$	attenuation coefficient
$L$	cylinder length
$m$	wavelength
$m_w$	cylinder mass
$MF_h$	added-mass of the water
$MF_T$	total horizontal longitudinal force(experimental)
$M_{S_1}$	total horizontal transverse force(experimental)
$M_{S_1}$	total longitudinal moment about any point $S_1$ on the cylinder
$MM_{S_1}$	$M_{S_1}$ (experimental)
$M_{ST}$	total transverse moment about any point $S_1$ on the cylinder
$MM_{ST}$	$M_{ST}$ (experimental)
$Re$	Reynolds number, $u_{max} D/\nu$
$S$	Strouhal number, $f_v D/u_{max}$
$t$	time
$T$	wave period
$u$	horizontal component of the water particle velocity

$u_{\max}$	maximum horizontal water particle velocity
$u_{\max} TAD$	Kuilegan-Carpenter number
$v$	vertical component of the water particle velocity
$x$	horizontal distance from the wave crest
$X$	horizontal distance from mean position of the wave generator
$y$	vertical distance from the still water level to the water particle
$\beta_{hf}$	phase angle of the maximum total horizontal longitudinal force
$\beta_{hm}$	phase angle of the maximum total longitudinal moment
$\beta_{ht}$	phase angle of the maximum total horizontal transverse force
$\delta$	unknown phase angle
$\epsilon_r$	reflection coefficient
$\mu$	Poisson's ratio
$\nu$	kinematic viscosity of fluid
$\rho$	mass density of water
$\sigma$	circular wave frequency, $2\pi/T$
$\partial u / \partial t$	horizontal component of the water particle local acceleration

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## 1. INTRODUCTION

This thesis consists of two parts; one deals with the calibration of a 200 ft x 15 ft x 10 ft (61 m x 4.57m x 3 m) wave tank, the other deals with wave forces on a vertical circular cylinder due to sinusoidal waves.

Experiments designed to verify the theoretical prediction of wave heights in terms of the stroke of the wave generator, water depth, and wave period were performed at Laboratoire Neyric (1952). Waves were generated by a paddle-type wave generator and the measured wave height at some distance from the wave generator was consistently about 30% below theoretical predictions. Ursell et al.<sup>27</sup> verified the linear wave generator theory by using a piston-type wave generator for a series of tests. When the wave steepness was within the range of  $0.002 \leq H/L \leq 0.03$ , the measured wave heights were 3.4% below the theoretical predictions. The measured wave heights were 10% below the theoretical prediction when the range of wave steepness was  $0.045 \leq H/L \leq 0.048$ . Therefore, the experiments verified the validity of the small-amplitude wave theory.

Madsen, Tenney, Keating and Webber<sup>13</sup> also studied waves using a piston-type wave generator. Madsen's results gave measured wave heights that were 15% below the theory. Tenney's experiments, in which two holes were drilled through the waveboard, indicated a difference of 3% between the two cases (holes open and closed). Keating and Webber measured wave heights in the range of  $0.005 \leq H/L \leq 0.08$  which were 6% below theoretical predictions if leakage around the wave generator could be eliminated. They concluded that linear wave generator theory was more applicable to shallow water waves than deep water waves.

The present experiments show good agreement, in the range of  $0.0012 \leq H/L \leq 0.0876$ , with linear wave generator theory for a piston-type wave generator.

Horizontal wave forces on a vertical circular cylinder which extends from the ocean bottom and pierces through the water surface are made up of two components. One is the longitudinal force in the direction of wave propagation, the other is the transverse force normal to the direction of wave propagation.

If the diameter of the cylinder is small compared with the wavelength ( $D/L < 0.1$ ), the longitudinal force is normally calculated by using Morison's equation which consists of a drag force and an inertia force. The transverse force due to vortex shedding from the cylinder, under certain circumstances, is not negligible. Wiegel, et al.<sup>29</sup> studied wave forces on piles and showed that large lateral oscillations could occur.

Jen<sup>12</sup> studied wave forces due to uniform periodic waves on the basis of linear wave theory in the ranges of  $d/L > 0.175$ ,  $H/L < 0.02$ , and  $D/L \leq 0.12$ . He found out that inertia forces were predominant and that the coefficient  $C_M$  had an overall average of 2.04. Morison et al.<sup>21</sup> presented data using small amplitude wave theory for  $0.102 < d/L < 0.529$ ,  $0.009 < H/L < 0.1135$ , and  $0.009 < D/L < 0.0419$ . He found out that  $C_M = 1.508 \pm 0.197$  and that  $C_D = 1.626 \pm 0.414$  (for  $2.2 \times 10^3 \leq Re \leq 1.11 \times 10^4$ ).

The experimental values of  $C_D$  and  $C_M$  were computed from the measured total moment in order to calculate the measured force, theoretical moment and force. In Jen's experiment, the drag forces on the circular pile were

so small that he assumed  $C_D$  to be zero and measured  $C_M$ . Wiegel et al.<sup>29</sup> calculated  $C_D$  when the acceleration was zero and  $C_M$  when the velocity was zero.

Wiegel et al.<sup>29</sup> showed that there was no relationship between  $C_M$  and Re. Morrison et al.<sup>21</sup> presented results showing that  $C_M$  and  $C_D$  were not dependent upon Re (for  $2.2 \times 10^3 < Re < 1,11 \times 10^4$ ). Wiegel<sup>28</sup> thought that the Keulegan-Carpenter number in terms of  $C_D$  and  $C_M$  was more significant than Re. Keulegan and Carpenter<sup>7</sup> stated that  $C_D$  and  $C_M$  were functions not of Re but of the Keulegan-Carpenter number. More recently it has been reported that both  $C_D$  and  $C_M$  are dependent on Re and Keulegan-Carpenter number<sup>2,7</sup>.

Bidde<sup>1</sup> showed that the transverse force was not dependent on Re but rather was a function of the Keulegan-Carpenter number. However, Isaacson and Maul<sup>10</sup> stated that the transverse forces were functions of Re and Keulegan-Carpenter number, and transverse forces occur when the Keulegan-Carpenter number had a value of 5 on the surface of the flow.

Jen<sup>12</sup>, Mogridge and Jamieson<sup>20</sup>, and Bidde<sup>1</sup> all report results of experimental investigations into wave forces on cylinders. They used a vertical circular cylinder in which the lower end of the cylinder was not in contact with the bottom of the tank. The present experiment, in which the cylinder is fixed on the bottom of the tank and wave forces and moments are measured over several ranges of wave parameters, yields reasonably good agreement with linear wave theory.

## 2. WAVE TANK CALIBRATION

### 2.1 Wave Generating System

The wave tank is a reinforced concrete structure which is 200 ft long, 15 ft wide and 10 ft deep, and has glass panels on one side near the test section, as shown in Figure 1. It is provided with a MTS closed-loop, servo-hydraulic wave generator (see Figure 2) which is capable of generating sinusoidal or programmed waves. Waves are produced by the translatory motion of a waveboard which is driven by the hydraulic actuator, as shown in Figure 3. The servo-valve regulates actuator stroke in response to the control signal developed in the servo controller. The servo controller develops the error signal by comparing its command and feedback. The waveboard position is automatically maintained at the command level. A block diagram of a position control servo loop is shown in Figure 4.

A beach of slope 1:6.2 (around  $9^{\circ}$ ), as shown in Figure 5, is constructed of 5 modules each consisting of orthogonal wooden grids mounted on top of plywood. Each module may be varied in elevation and slope. Figure 6 shows the layout of the wave tank.

Each wave probe is of the resistance type consisting of a pair of parallel stainless steel wires, 0.059 in. in diameter, 19.69 in. in length and spaced 0.49 in. apart. The wave probes are calibrated in still water by raising and lowering them by known increments. The output of each probe was recorded on magnetic tape and on a strip chart recorder.

### 2.2 Wave Generator Performance

The closed loop servo-hydraulic wave generator is designed to generate wave forms within the wave height to wavelength performance envelope, at different water depth, on the basis of Biesel's equation

for a piston-type wave generator. Biesel's equation (linear wave generator theory) is defined as:

$$H = \frac{e \sinh^2 \left( \frac{2\pi d}{L} \right)}{\sinh \left( \frac{2\pi d}{L} \right) \cosh \left( \frac{2\pi d}{L} \right) + \frac{2\pi d}{L}} \quad (1)$$

where  $e$  is the stroke of the wave generator.

Two wave probes were installed about 35 ft. apart near the middle of the wave tank to measure the wave height and wave length. The wave height to wave length performance envelope was then established by varying the stroke of the wave generator from 0.5 in to 10 in, at frequencies of 0.2 Hz to 1 Hz with 4.5 ft and 3 ft water depths as shown in Figures 7 and 9. The computer program based on the Biesel's equation for the analysis of the piston-type wave generator is listed in Appendix A.

### 2.3 Reflection Coefficient of the Beach

When a primary incident wave travels toward the beach, it will be partially reflected and partially absorbed by the beach.

Successive reflections occur until steady state is established. A partial standing wave system is thus established along the length of the wave tank. Therefore, the wave heights are not the same at all points but vary as a function of the distance from the wave generator.<sup>27</sup>

$$H = 2a[1 + e_r \cos(2kx_1 + \delta) + e_r \cos \delta] \quad (2)$$

That is, the wave height variation is a sinusoidal oscillation with a wavelength which is approximately equal to one-half the wavelength of the incident wave.

The definition of the reflection coefficient  $e_r$  of the beach is the ratio of the reflected wave height ( $H_r$ ) to the incident wave height ( $H_i$ ).

If the maximum and minimum values of the sinusoidal variation in wave height are denoted by  $H_{\max}$  and  $H_{\min}$ , respectively, then the reflection coefficient may be expressed as:

$$r = \frac{H_r}{H_i} = \frac{H_{\max} - H_{\min}}{H_{\max} + H_{\min}} \quad (3)$$

Seven wave probes spaced 2 ft. apart, the first being located 48 ft from the mean position of the wave generator in the wave tank were used to measure reflection coefficients. (see Fig. 13).

#### 2.4 Wave Attenuation Coefficient

When wave heights are measured in a finite width wave tank, the viscosity of the boundary layer on the side walls and bottom of the tank causes a slight attenuation of the wave height due to energy dissipation in the wave propagation. A theoretical estimation of the attenuation coefficient  $K$  is defined as

$$K = \frac{2k}{b} \left( \frac{v}{2\sigma} \right)^{1/2} \frac{(kb + \sinh 2kd)}{2kd + \sinh 2kd} \quad (4)$$

where  $b$  is the width of the wave tank,  $v$  is the kinematic viscosity,  $k$  is the wave number and  $\sigma$  is  $\frac{2\pi}{T}$ .

Suppose that two wave probes are located 61.8 ft apart near the middle of the wave tank to avoid the wave generator and beach effects and the individual wave heights  $H_1$  and  $H_2$  are measured. Then the attenuation coefficient can be determined by:

$$H_2 = H_1 e^{-K(x_2 - x_1)} \quad (5)$$

where  $x_2 - x_1 = 61.8$  ft.

Obviously, the decrease in the wave height is so small that it can be neglected from a practical point of view. The theoretical and measured wave attenuation coefficient are shown in Table 3.

## 2.5 Discussion of Results

A series of tests were performed in order to verify linear wave generator theory. This theory predicts the relationship between the wave height, wave period, water depth and the stroke of the wave generator for a piston-type wave generator.

Measured values of  $H$  and  $L$ ,  $H/e$  and  $2\pi d/L$ , for 4.5 ft and 3 ft water depths over the range of  $0.074 < d/L < 0.6977$ ,  $0.0012 < H/L < 0.0876$  and  $0.058 < d/L < 0.548$ ,  $0.0012 < H/L < 0.1$  respectively are compared with those obtained from equation (1) in Figures 7-12 and Tables 1 and 2. The wavelength might be considered as a constant throughout the entire range of strokes of the wave generator at the same wave period as it has a maximum standard deviation of 1.38 and 0.58 for 4.5 ft and 3 ft water depths respectively. The wave generator can generate a maximum wave height of 1.343 ft having a corresponding wavelength of 13.35 ft. The maximum wave steepness is 1/10. The wave height to wavelength performance envelopes for the 3 ft water depth, as shown in Figures 9 and 10, have a better agreement with linear wave generator theory than the wave height to wavelength performance envelopes for the 4.5 ft water depth, as shown in Figures 7-8.

The experimental results in the range of the small amplitude waves  $0.0012 < H/L < 0.03$  (Figures 11 and 12) show better agreement with theory than in the range of the finite amplitude waves  $0.03 < H/L < 0.0876$  and  $0.03 < H/L < 0.1$  for both 4.5 ft and 3 ft water depths respectively.

However, all the experimental results for a 3 ft water depth are closer to the theory than those results obtained for the 4.5 ft water depth. A variety of factors could account for the difference. Errors in measurement of wave height, wavelength, water depth, and wave generator stroke; some effects such as interference due to transverse reflection from the side walls of the wave tank; leakage around the waveboard; motion of the waveboard which is not quite simple harmonic; and lack of a perfect wave energy absorber.

The wave height variation with distance along the wave tank is a sinusoidal oscillation which has a wavelength approximately equal to one-half the wavelength of the incident wave. The reflection coefficient is less than 10%. The attenuation coefficient of the wave height due to viscous boundary layers is so small that it can be neglected from a practical standpoint, as shown in Table 3.

### 3. RESPONSE OF A VERTICAL CIRCULAR CYLINDER TO SINUSOIDAL WAVES

#### 3.1 Wave Induced Forces and Overturning Moments

Morison's equation for the wave forces due to potential flow on a submerged vertical cylinder for which the ratio of cylinder diameter to wavelength is very small consists of two components. One is a drag force due to the water particle velocity, the other is an inertia force due to the water particle acceleration. These coefficients may be determined experimentally. It is assumed that the flow is inviscid, incompressible, and irrotational. Furthermore, the ratio of wave height to wave length is so small that linear wave theory can be applied. The inertia force on the submerged vertical circular cylinder may be expressed as:

$$F_I = C_{M0} \left( \frac{\pi D^2}{4} \right) \frac{\partial u}{\partial t} \quad (6)$$

The drag force which consists of both frictional drag and form drag on the submerged vertical circular cylinder may be expressed as:

$$F_D = \frac{1}{2} C_D \rho A u^2 \quad (7)$$

Morison, O'Brien and Johnson (1950) combined the inertia force and drag force to obtain the total force expressed as:

$$F = F_I + F_D = C_{M0} \left( \frac{\pi D^2}{4} \right) \frac{\partial u}{\partial t} + \frac{1}{2} C_D \rho A |u| u \quad (8)$$

where the horizontal components of velocity and local acceleration are given by:

$$u = \frac{\pi H}{T} \frac{\cosh[2\pi(y+d)/L]}{\sinh 2\pi d/L} \cos 2\pi \left( \frac{x}{L} - \frac{t}{T} \right) \quad (9)$$

$$\frac{\partial u}{\partial t} = \frac{2\pi^2 H}{T^2} \frac{\cosh[2\pi(y+d)/L]}{\sinh 2\pi d/L} \sin 2\pi \left(\frac{x}{L} - \frac{t}{T}\right) \quad (10)$$

where  $x, t = 0$  defines the wave crest to be at the centre of the cylinder.

The horizontal force  $F_h(s)$  on a differential length of a vertical circular cylinder  $ds$ , as shown in Figure 13, is given by

$$F_h(s) = [C_M p (\frac{\pi D}{4})] \frac{\partial u}{\partial t} + \frac{1}{2} C_D p D |u| u ds \quad (11)$$

It is assumed that  $C_D$  and  $C_M$  are constant for the entire length of the cylinder. The total horizontal force on the cylinder between two points  $S_1$  and  $S_2$  where  $S_2 < S_3$  (the surface) is given by:

$$F_h = \int_{S_1}^{S_2} F_h(s) ds = \int_{S_1}^{S_2} [C_M p (\frac{\pi D}{4})] \frac{\partial u}{\partial t} + \frac{1}{2} C_D p D |u| u ds \\ = \pi p D \frac{H^2 L}{T^2} \left[ \frac{\pi D}{4H} C_M K_2 \sin 2\pi \left(\frac{x}{L} - \frac{t}{T}\right) + C_D K_1 [\cos 2\pi \left(\frac{x}{L} - \frac{t}{T}\right) \right. \\ \left. - \cos 2\pi \left(\frac{x}{L} - \frac{t}{T}\right)] \right] \quad (12)$$

$$\text{where } K_1 = \frac{\frac{4\pi S_2}{L} - \frac{4\pi S_1}{L} + \sinh \frac{4\pi S_2}{L} - \sinh \frac{4\pi S_1}{L}}{16 (\sinh \frac{2\pi d}{L})^2} \quad (13)$$

$$K_2 = \frac{\sinh \frac{2\pi S_2}{L} - \sinh \frac{2\pi S_1}{L}}{\sinh \frac{2\pi d}{L}} \quad (14)$$

By differentiating equation (12) with respect to  $2\pi \left(\frac{x}{L} - \frac{t}{T}\right)$  and

setting the derivative  $dF_h \times 2\pi(x/L-t/T)$  equal to zero, the phase angle  $\beta_{hf}$  of the maximum total horizontal force can be obtained.

$$\sin \beta_{hf} = + \frac{\pi D C_M K_2}{8 H C_D K_1} \quad (15)$$

If  $\sin \beta_{hf} > 1$ , then  $\cos \beta_{hf} = 0$ ; that is, the wave crest (or trough) lags behind the force by  $\pi/2$ . (see Ref. 28)

The total longitudinal moment about any point  $S_1$  on the vertical circular cylinder is given:

$$\begin{aligned} M_{S_1} &= \int_{S_1}^{S_2} (S - S_1) dF_h \\ &= \rho D \left(\frac{H L}{T}\right)^2 \left\{ \frac{\pi D}{4 H} C_M K_4 \sin 2\pi \left(\frac{x}{L} - \frac{t}{T}\right) + C_D K_3 |\cos 2\pi \left(\frac{x}{L} - \frac{t}{T}\right)| \cos 2\pi \left(\frac{x}{L} - \frac{t}{T}\right) \right. \\ &\quad \left. - \frac{2\pi S_1}{L} \left[ \frac{\pi D}{8 H} C_M K_2 \sin 2\pi \left(\frac{x}{L} - \frac{t}{T}\right) + \frac{1}{2} C_D K_1 |\cos 2\pi \left(\frac{x}{L} - \frac{t}{T}\right)| \cos 2\pi \left(\frac{x}{L} - \frac{t}{T}\right) \right] \right\} \end{aligned} \quad (16)$$

where  $K_1$  and  $K_2$  are same as equations (13) and (14).

$$K_3 = \frac{\frac{1}{2} \left( \frac{4\pi S_2}{L} \right)^2 - \frac{1}{2} \left( \frac{4\pi S_1}{L} \right)^2 + \frac{4\pi S_2}{L} \sinh \frac{4\pi S_2}{L} - \frac{4\pi S_1}{L} \sinh \frac{4\pi S_1}{L}}{64 \left( \sinh \frac{2\pi d}{L} \right)^2} - \cosh \frac{4\pi S_2}{L} + \cosh \frac{4\pi S_1}{L} \quad (17)$$

$$K_4 = \frac{\frac{2\pi S_2}{L} \sinh \frac{2\pi S_2}{L} - \frac{2\pi S_1}{L} \sinh \frac{2\pi S_1}{L} - \cosh \frac{2\pi S_2}{L} + \cosh \frac{2\pi S_1}{L}}{2 \sinh \left( \frac{2\pi d}{L} \right)} \quad (18)$$

By differentiating equation (16) with respect to  $2\pi \left( \frac{x}{L} - \frac{t}{T} \right)$  and setting the derivative  $dM_{S_1}/d\{2\pi(x/L-t/T)\}$  equal to zero, the phase angle  $\beta_{hm}$  of the maximum total longitudinal moment at  $S_1$  can be obtained.

$$\sin \beta_{hm} = \pm \frac{\frac{2\pi S_1}{L} K_2}{8 H C_D [K_3 - \left( \frac{2\pi S_1}{L} \right) \left( \frac{K_1}{2} \right)]} \quad (19)$$

If  $\sin \beta_{hm} > 1$ , then  $\cos \beta_{hm} = 0$ ; that is, the wave crest (or trough) lags behind the moment by  $\pi/2$ . (see Ref. 28)

The transverse force, due to vortex shedding as the waves pass the vertical circular cylinder, may be expressed as

$$F_L = \frac{1}{2} C_L \rho u^2 A \quad (20)$$

and the transverse force  $F_T(s)$  on a differential length of a vertical circular cylinder  $ds$  is given

$$F_T(s) = \left[ \frac{1}{2} C_L \rho u^2 D \right] ds \quad (21)$$

Assuming that  $C_L$  is a constant for the entire length of the cylinder, the total transverse force on the cylinder between  $S_1$  and  $S_2$  is given

$$F_T = \int_{S_1}^{S_2} F_T(s) ds = \int_{S_1}^{S_2} \frac{1}{2} C_L \rho u^2 D ds = \frac{C_L \rho D \pi H^2 L K_1}{T^2} \{\cos 2\pi(\frac{x}{L} - \frac{t}{T})\}^2 \quad (22)$$

where  $K_1$  is same as equation (13).

By differentiating equation (22) with respect to  $2\pi(x/L - t/T)$  and setting the derivative  $dF_T/d[2\pi(x/L - t/T)]$  equal to zero, the phase angle  $\beta_{hT}$  of the maximum total horizontal transverse force can be obtained.

$$\sin 2\beta_{hT} = 0 \quad (23)$$

therefore  $\beta_{hT} = 0^\circ$  or  $\pi/2$ .

$$\text{Therefore, } F_{T\max} = \frac{C_L \rho D \pi H^2 L K_1}{T^2} \quad (24)$$

$$\text{or } F_{T\min} = 0$$

The total transverse moment about any point  $S_1$  on the vertical circular cylinder is given

$$M_{ST} = \int_{S_1}^{S_2} (S - S_1) dF_T \\ = C_L \rho D \left(\frac{H}{T}\right)^2 \left(K_3 - \frac{\pi S_1}{L} K_1\right) \{\cos 2\pi(\frac{x}{L} - \frac{t}{T})\}^2 \quad (25)$$

where  $K_1$  and  $K_3$  are same as equations (13) and (17).

By differentiating equation (25) with respect to  $2\pi(x/L - t/T)$  and setting the derivative  $dM_{ST}/d[2\pi(x/L - t/T)]$  equal to zero, the phase angle  $\beta_{MST}$  of the maximum total transverse moment at  $S_1$  can be obtained.

$$\sin 2\beta_{MST} = 0 \quad (26)$$

therefore  $\beta_{MST} = 0^\circ$  or  $\pi/2$ .

Therefore, when  $\beta_{MST} = \pi/2$ ,  $M_{ST} = 0$  is the minimum value; when  $\beta_{MST} = 0^\circ$ ,

$$M_{ST} = C_L \rho D \left( \frac{H}{T} \right)^2 \left( K_3 - \frac{2\pi S_1}{L} \frac{K_1}{2} \right) \quad (27)$$

### 3.2 Keulegan-Carpenter Number and Reynolds Number

There are two nondimensional numbers, Reynolds number and Keulegan-Carpenter number, which are related to wave forces and drag, inertia, and lift coefficients.  $C_D$  is a function of  $Re$  in steady flow. Hogben et al.<sup>7</sup> have shown that  $C_D$  decreased with increasing  $Re$  in steady flow (for  $5 \times 10^4 \leq Re \leq 5 \times 10^5$ ). For a small diameter cylinder  $C_D$ ,  $C_M$ , and  $C_L$  depend on  $u_{max} T/D$ . Keulegan and Carpenter have shown that  $C_D$  and  $C_M$  of the circular cylinder were functions of  $u_{max} T/D$ . When  $u_{max} T/D$  is small, separation of flow from the cylinder doesn't occur; when  $u_{max} T/D$  is relatively high, a Karman vortex street is formed.

The definitions of Keulegan-Carpenter number and Reynolds number may be expressed as

$$\text{Keulegan-Carpenter number} = \frac{u_{max} T}{D} = \frac{\pi H \cosh \{2\pi(y + d)/L\}}{D \sinh 2\pi d/L} \quad (28)$$

$$\text{Reynolds number} = \frac{u_{max} D}{v} = \frac{\pi H D \cosh \{2\pi(y + d)/L\}}{T v \sinh 2\pi d/L} \quad (29)$$

### 3.3 Wave-Induced Vibration of a Vertical Circular Cylinder

The natural frequencies of a circular cylinder as shown in Figure 13 are given by

$$f_i = \frac{a_i}{2\pi} \left[ \frac{EI}{(m + m_w) \ell^3} \right]^{1/2} \quad (30)$$

where  $I = \pi D^4/64$  is the moment of inertia, E is the Elastic Modulus,  $\ell$  is the cylinder length,  $m$  is the mass of the cylinder,  $m_w$  is the added-mass of the water. Added-mass per unit length for a circular cross section is equal to  $\rho \pi r^2$ , where  $r$  is the radius of cylinder, and  $\rho$  is the mass density of water. Crede (1965) pointed out that the first three values of modal coefficient,  $a_i$ , are 3.52, 22.4 and 61.7.

The designer can adjust the values of E, I, and  $\ell$  in order to make sure that the natural frequencies of the circular cylinder do not fall within the range of surface wave frequencies. A resonance condition may thus be avoided.

When waves pass a vertical circular cylinder, lateral vibration occurs due to vortex shedding. Under certain circumstances, this lateral vibration is not negligible. The destruction of the Texas Tower No. 4 off the coast of New Jersey (January 15, 1961) is attributed to coupling between the natural frequency of the structure and the vortex shedding frequency.

The non-dimensional Strouhal number S is used to define the vortex shedding frequency  $f_v$  which is a function of the cylinder diameter D and the maximum horizontal particle velocity  $u_{max}$ , where

$$S = \frac{f_v D}{u_{max}} \quad (31)$$

For circular cylinders, the empirical average value of the Strouhal number is about 0.2 for a sub-critical Reynolds number  $Re \leq 2 \times 10^5$  (Ref. 28, Delany and Sorensen, 1953; Fung, 1960; Humphreys, 1960) in uniform rectilinear flow.

### 3.4 Froude's Model Laws

Froude's scaling laws are applied to hydraulic models when viscous and capillary effects are negligible. Froude's scaling system may be described as follows:

$$\text{Length Scale: } \frac{\text{Length of model}}{\text{Length of prototype}} = n_L$$

$$\text{Time Scale: } n_t = n_L^{1/2}$$

$$\text{Velocity Scale: } n_v = n_L^{1/2}$$

$$\text{Mass Scale: } n_m = n_L^3$$

$$\text{Force Scale: } n_F = n_L^3$$

$$\text{Pressure Scale: } n_p = n_L$$

$$\text{Shearing Stress Scale: } n_t = n_L$$

$$\text{Drag Coefficient Scale: } n_{CD} = 1$$

$$\text{Mass Coefficient Scale: } n_{CM} = 1$$

$$\text{Elastic Modulus Scale: } n_E = n_L$$

$$\text{Poisson's Ratio Scale: } n_\mu = 1$$

### 3.5 Experimental Technique

A rigid Poly Vinyl Chloride circular cylinder was used for all experiments. It had a length of 5 ft, 6 in. inside diameter and 9/32 in. wall thickness. The cylinder was attached to a 1 in. thick circular

aluminum base which was bolted to a 2 ft. diameter by 1/2 in. thick steel plate. This steel plate was anchor bolted to the bottom of the wave tank and a wooden annulus, ground board, was bolted to the steel plate.

Eight 120 ohm strain gauges were mounted vertically around the outside of the cylinder near the steel base, and were waterproofed as shown in Figure 15. The technique for measuring total moment and force was based on elastic beam theory. One Wheatstone bridge, made up of four active strain gauges to produce an output due to the tension or compression of the cylinder, was used to measure the longitudinal overturning moment about the base of the cylinder. Another Wheatstone bridge was used to measure the transverse overturning moment.

The cylinder was calibrated by applying loads to different points along the length of the cylinder. These loads varied from 2 lbs to 20 lbs, as shown in Figure 15 and Table 4. The calibration indicated a linear relationship between the loads applied and the strain readings. Therefore the values of strain obtained could be interpreted as the longitudinal or transverse moments about a reference height,  $S_1$ , (see Figure 13).

The centre of the cylinder was located at the centre line of the wave tank and 48 ft. from the static position of the wave generator. A wave gauge was located alongside the centre of the cylinder, midway between the cylinder and the tank wall to measure the wave heights, which were approximately the same as the incident wave heights measured without the cylinder in the wave tank, as shown in Figure 16.

Sinusoidal waves were generated in 3 ft. water depth throughout the tests. Four different wave periods 1.429 sec, 1.667 sec, 2 sec, and 2.5 sec were run and ten different wave heights were generated for each wave period. Data was recorded after the waves passing the cylinder reached a steady-state condition. A digital strain indicator (B and K type, 1526) and a two-channel Hewlett-Packard 71003 strip chart recorder were used for this recording system, as shown in Figure 18.

The value of Elastic Modulus was found experimentally to be  $5.073 \times 10^5$  psi. Figure 19 shows the tension specimen mounted in an Instron testing machine.

#### DISCUSSION OF RESULTS

The solid line represents theory in all graphs. The values of  $C_D$ ,  $C_M$ , and  $C_L$  must be determined experimentally. When the crest (or trough) passes the center line of cylinder (zero acceleration),  $2\pi(\frac{x}{L} - \frac{t}{T}) = 0^\circ$ ,  $C_D$  and  $C_L$  may be solved from equations (16) and (25). When the instantaneous water surface profile is at the still water level (zero velocity),  $2\pi(\frac{x}{L} - \frac{t}{T}) = 90^\circ$ ,  $C_M$  may be solved from equation (16). Then the experimentally determined values of  $C_D$ ,  $C_M$ , and  $C_L$  are used for computing theoretical moment and force. The action line of total resultant force to the cylinder may be calculated and a static calibration (TABLE 4) is utilized for computing the measured force.

The computer program based on linear wave theory which estimates the wave forces and overturning moments on the circular cylinder is given

in Appendix B. Values of  $C_D$ ,  $C_M$ , and  $C_L$  may be obtained from this program.

Results obtained from the program are given in Table 5.

The drag, mass, and lift coefficients are functions of Reynolds and Keulegan-Carpenter number (see Figures 20-25). When the ranges of Reynolds and Keulegan-Carpenter numbers are between  $8.286 \times 10^3$  and  $2.083 \times 10^5$  (subcritical regime), 0.6923 and 11.83, respectively, values of  $C_M$  and  $C_L$  vary between 0.675 to 1.22 and 0 to 0.17 respectively.  $C_M$  and  $C_L$  might be considered as constant for the wave periods 1.429 sec., 1.667 sec., and 2 sec. However,  $C_D$  varies from 0.307 to 7.418, as shown in Figures 20 and 21: In addition  $C_D$  is inversely related to  $H$  and  $u_{max}$ . That is, when the wave height becomes very small,  $C_D$  becomes very large, and  $u_{max}$  approaches zero. The wide variation of  $C_D$  has also been observed by a number of other investigators. Jan<sup>12</sup>, and Pratte et al.<sup>24</sup> have presented data for the wave forces on a 6 in. diameter pile in 3 ft. water depth based on linear wave theory.

Values of  $C_D$  and  $C_M$  are shown in Table 5 to be independent of  $d/L$ ,  $D/L$ , and  $d/T^2$ . Figures 26-29 show the maximum longitudinal and transverse forces versus Reynolds number and Keulegan-Carpenter number for four different wave periods. The longitudinal and transverse forces increase with increasing  $Re$  and  $u_{max} T/D$ . The transverse forces occur when the Keulegan-Carpenter number at the surface is greater than 5.27. Bidde<sup>1</sup> reported values of 3 to 5. Keulegan and Carpenter have observed a value of 15 for a horizontal rather than a vertical cylinder, using a standing wave with  $d/L$  being very small. When the Reynolds number is greater than  $7.397 \times 10^4$ , eddy shedding occurs.

Maximum forces and overturning moments are plotted as functions of wave steepness  $H/L$  for four different wave periods. The range of  $d/L$  is between 0.123 and 0.277, the range of  $H/L$  is between 0.003 and 0.105, and the range of  $D/L$  is between 0.023 and 0.051. Figures 30-33 show that the maximum forces and overturning moments are increasing functions of  $H/L$ . The theoretical and experimental results show reasonably good agreement. It should be noted that the maximum values of moment were taken as the mean value of all maxima observed.

Figure 17 shows symmetric vortex shedding behind the cylinder. The fundamental frequency of the cylinder is much higher than the wave frequency and the vortex-shedding frequency (equations (30) and (31)). Therefore, resonance does not occur.

#### 4. CONCLUSIONS

The piston-type wave generator is capable of generating deep, intermediate and shallow water waves. A series of experiments have verified the ability of linear wave generator theory to predict wave motion associated with the stroke of the wave generator, wave height, water depth, and wave period. The theoretical and experimental results show better agreement for shallow water waves than intermediate water waves.

The measured wave forces and overturning moments on small diameter circular cylinder based on the linear wave theory show reasonably good agreement with theoretical results for the ranges of  $0.123 < d/L < 0.277$ ,  $0.003 < H/L < 0.105$ , and  $0.023 < D/L < 0.051$  in the wave periods of 1.429 sec., 1.667 sec., 2 sec., and 2.5 sec. The coefficients  $C_M$  and  $C_L$  have been found to have an overall average of 0.953 and 0.0535 respectively and might be considered as constants; while  $C_D$  was found to vary between 7.418 and 0.307 when the Reynolds number was in the range of  $8.286 \times 10^3$  to  $2.083 \times 10^5$  and the Keulegan-Carpenter number varied between 0.6923 and 11.83.

Transverse forces are related to vortex shedding. When the Keulegan-Carpenter number at the surface is greater than 5.27, transverse forces occur. Although the transverse forces and corresponding moments are small, they are not negligible when scaled up to prototype values (Froude's similitude requires that the force scale be equal to the third power of the length scale).

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TABLE 1

Calibration Results: Wave Parameters for 4.5 ft Water Depth

f(Hz)	Span(%)	e(in.)	H(ft)	L(ft)	d/L	H/L	$2\pi d/L$	H/e
0.2	10	0.94	0.07	59.5	0.076	0.0012	0.49	0.90
0.2	20	1.88	0.16	60.0	0.075	0.0027	0.49	1.03
0.2	30	2.88	0.31	56.9	0.079	0.0055	0.49	1.29
0.2	40	3.88	0.34	56.0	0.080	0.0061	0.49	1.05
0.2	50	4.88	0.39	55.8	0.080	0.0070	0.49	0.96
0.2	60	5.88	0.56	56.5	0.080	0.0099	0.49	1.14
0.2	70	6.75	0.56	60.7	0.074	0.0092	0.49	1.00
0.2	80	7.63	0.64	56.0	0.080	0.0110	0.49	1.00
0.2	90	8.69	0.75	56.7	0.079	0.0130	0.49	1.53
<hr/>								
0.3	10	0.78	0.10	36.69	0.123	0.0027	0.72	1.54
0.3	20	1.59	0.19	37.26	0.121	0.0051	0.72	1.43
0.3	30	2.44	0.23	41.25	0.109	0.0056	0.72	1.13
0.3	40	3.28	0.22	38.50	0.117	0.0057	0.72	1.54
0.3	50	4.13	0.56	39.75	0.113	0.0141	0.72	1.63
0.3	60	4.88	0.70	39.10	0.115	0.0179	0.72	1.75
0.3	70	5.56	0.81	41.25	0.109	0.0196	0.72	1.76
0.3	80	6.41	0.81	41.25	0.109	0.0196	0.72	1.53
0.3	90	7.59	1.01	38.50	0.117	0.0262	0.72	2.53
<hr/>								
0.4	10	0.63	0.22	25.74	0.175	0.0085	1.10	4.23
0.4	20	1.34	0.23	25.74	0.175	0.0089	1.10	2.05
0.4	30	2.06	0.34	25.74	0.175	0.0132	1.10	1.98
0.4	40	2.75	0.47	25.74	0.175	0.0183	1.10	2.05
0.4	50	3.06	0.59	25.74	0.175	0.0229	1.10	2.02
0.4	60	4.09	0.68	25.74	0.175	0.0264	1.10	2.00
0.4	70	5.56	0.80	25.74	0.175	0.0311	1.10	1.74
0.4	80	5.38	0.96	25.74	0.175	0.0373	1.10	2.13

TABLE 1-Continued

f(Hz)	Span(%)	e(in.)	H(ft)	L(ft)	d/L	H/L	$2\pi d/L$	H/e
0.5	10	0.53	0.13	19.18	0.2346	0.0060	1.43	2.95
0.5	20	1.13	0.31	19.66	0.2289	0.0158	1.43	3.33
0.5	30	1.75	0.47	19.86	0.2266	0.0237	1.43	2.73
0.5	40	2.38	0.65	20.00	0.2250	0.0325	1.43	3.28
0.5	50	2.56	0.71	20.00	0.2250	0.0355	1.43	3.32
0.5	60	3.53	0.94	20.58	0.2186	0.0457	1.43	3.24
0.5	70	4.09	1.07	19.44	0.2315	0.0550	1.43	3.15
0.5	80	4.63	1.19	19.93	0.2258	0.0597	1.43	3.05
0.5	90	5.06	1.28	20.00	0.2250	0.0640	1.43	4.44
0.6	10	0.89	0.13	15.31	0.2939	0.0085	1.89	3.61
0.6	20	0.97	0.33	14.69	0.3063	0.0225	1.89	3.83
0.6	30	1.53	0.47	14.50	0.3103	0.0324	1.89	3.67
0.6	40	2.06	0.61	14.89	0.3022	0.0410	1.89	3.55
0.6	50	2.25	0.72	14.89	0.3022	0.0484	1.89	3.83
0.6	60	3.03	0.87	14.89	0.3022	0.0584	1.89	3.48
0.6	70	3.59	0.94	14.88	0.3022	0.0632	1.89	3.13
0.6	80	4.03	1.02	15.00	0.3000	0.0680	1.89	3.00
0.7	10	0.81	0.17	11.03	0.4080	0.0154	2.50	5.00
0.7	20	0.84	0.34	11.03	0.4080	0.0308	2.50	4.86
0.7	30	1.34	0.50	10.89	0.4132	0.0459	2.50	4.46
0.7	40	1.81	0.59	11.88	0.3788	0.0497	2.50	4.37
0.7	50	2.03	0.74	11.14	0.4039	0.0664	2.50	4.38
0.7	60	2.72	0.82	11.14	0.4039	0.0736	2.50	3.57
0.7	70	3.19	0.96	11.53	0.3903	0.0833	2.50	3.56
0.7	80	3.59	0.99	11.53	0.3903	0.0859	2.50	3.30

TABLE 1-Continued

$f$ (Hz)	Span(%)	$e$ (in.)	$H$ (ft)	$L$ (ft)	$d/L$	$H/L$	$2\pi d/L$	$H/e$
0.8	10	0.34	0.15	8.57	0.5251	0.0175	3.30	5.17
0.8	20	0.72	0.47	8.27	0.5441	0.0568	3.30	7.83
0.8	30	1.18	0.55	8.17	0.5508	0.0673	3.30	5.56
0.8	40	1.63	0.60	7.96	0.5653	0.0754	3.30	4.44
0.8	50	1.81	0.69	8.93	0.5039	0.0773	3.30	4.64
0.8	60	2.44	0.79	9.02	0.4989	0.0876	3.30	3.95
0.9	10	0.36	0.12	6.45	0.6977	0.0186	4.21	4.61
0.9	20	0.60	0.24	6.49	0.6934	0.0370	4.21	4.53
0.9	30	1.08	0.35	6.50	0.6923	0.0538	4.21	3.85
0.9	40	1.44	0.52	6.88	0.6541	0.0756	4.21	4.26
0.9	50	1.68	0.58	7.07	0.6365	0.0820	4.21	4.30

TABLE 2

Calibration Results: Wave Parameters for 3 ft Water Depth

f(Hz)	Span(%)	e(in.)	H(ft)	L(ft)	d/L	H/L	$2\pi d/L$	H/e
0.2	10	0.97	0.06	50.00	0.060	0.0012	0.38	0.75
0.2	20	2.00	0.13	50.72	0.059	0.0026	0.38	0.76
0.2	30	3.00	0.18	51.47	0.058	0.0035	0.38	0.72
0.2	40	4.03	0.27	50.00	0.060	0.0054	0.38	0.79
0.2	50	5.09	0.34	50.72	0.059	0.0067	0.38	0.81
0.2	60	6.13	0.41	51.47	0.058	0.0080	0.38	0.80
0.2	70	7.13	0.49	50.00	0.060	0.0098	0.38	0.83
0.2	80	8.16	0.56	50.00	0.060	0.0112	0.38	0.82
0.2	90	9.06	0.67	50.72	0.059	0.0132	0.38	0.88
0.2	100	10.00	0.74	50.70	0.059	0.0146	0.38	0.89
0.3	10	0.78	0.085	33.06	0.09	0.0026	0.57	1.29
0.3	20	1.75	0.19	33.03	0.09	0.0058	0.57	1.27
0.3	30	2.63	0.27	33.50	0.09	0.0080	0.57	1.23
0.3	40	3.53	0.37	33.03	0.09	0.0110	0.57	1.28
0.3	50	4.38	0.47	33.50	0.09	0.0140	0.57	1.31
0.3	60	5.25	0.56	32.53	0.09	0.0170	0.57	1.27
0.3	70	6.09	0.67	33.00	0.09	0.0200	0.57	1.31
0.3	80	6.97	0.74	33.03	0.09	0.0220	0.57	1.28
0.3	90	7.81	0.82	33.50	0.09	0.0240	0.57	1.26
0.3	100	8.63	0.91	33.50	0.09	0.0270	0.57	1.26
0.4	10	0.68	0.08	23.33	0.1286	0.0034	0.80	1.33
0.4	20	1.44	0.21	23.33	0.1286	0.0090	0.80	1.91
0.4	30	2.25	0.34	23.65	0.1286	0.0144	0.80	1.79
0.4	40	3.06	0.48	23.03	0.1300	0.0208	0.80	1.85
0.4	50	3.81	0.58	24.30	0.1235	0.0239	0.80	1.81
0.4	60	4.53	0.68	23.03	0.1300	0.0295	0.80	1.79
0.4	70	5.25	0.80	22.73	0.1320	0.0352	0.80	1.82
0.4	80	5.97	0.89	23.65	0.1268	0.0376	0.80	1.78
0.4	90	6.66	0.98	24.31	0.1234	0.0400	0.80	1.78
0.4	100	7.31	1.06	24.31	0.1234	0.0436	0.80	1.74

TABLE 2-Continued

f(Hz)	Span(%)	e(in.)	H(ft)	L(ft)	d/L	H/L	$2\pi d/L$	H/e
0.5	10	0.59	0.12	17.72	0.169	0.0068	1.05	2.40
0.5	20	1.75	0.27	18.18	0.165	0.0149	1.05	1.80
0.5	30	1.94	0.42	17.95	0.167	0.0234	1.05	2.62
0.5	40	2.63	0.54	20.00	0.150	0.0270	1.05	2.45
0.5	50	3.28	0.69	17.95	0.167	0.0384	1.05	2.56
0.5	60	3.86	0.80	17.72	0.169	0.0450	1.05	2.50
0.5	70	4.53	0.92	17.95	0.167	0.0510	1.05	2.42
0.5	80	5.09	1.00	18.18	0.165	0.0550	1.05	2.38
0.5	90	5.69	1.14	17.95	0.167	0.0635	1.05	2.43
0.5	100	6.31	1.33	17.95	0.167	0.0740	1.05	2.51
0.6	10	0.50	0.13	13.13	0.2285	0.0099	1.41	3.25
0.6	20	1.06	0.28	13.43	0.2234	0.0210	1.41	3.11
0.6	30	1.69	0.45	13.68	0.2193	0.0329	1.41	3.21
0.6	40	2.31	0.58	13.58	0.2210	0.0427	1.41	3.05
0.6	50	2.88	0.74	13.13	0.2285	0.0560	1.41	3.08
0.6	60	3.38	0.85	13.18	0.2285	0.0650	1.41	3.04
0.6	70	3.88	0.96	13.58	0.2210	0.0700	1.41	3.00
0.6	80	4.44	1.20	13.18	0.2285	0.0910	1.41	3.24
0.6	90	4.94	1.27	13.35	0.2247	0.0948	1.41	3.08
0.6	100	5.50	1.34	13.35	0.2247	0.1006	1.41	2.93
0.7	10	0.47	0.14	10.82	0.277	0.0120	1.70	3.50
0.7	20	0.94	0.30	11.07	0.271	0.0271	1.70	3.75
0.7	30	1.50	0.44	11.05	0.271	0.0398	1.70	3.38
0.7	40	2.00	0.58	10.94	0.274	0.0530	1.70	3.41
0.7	50	2.50	0.71	10.94	0.274	0.0649	1.70	3.38
0.7	60	3.00	0.82	11.41	0.263	0.0719	1.70	3.28
0.7	70	3.50	0.90	11.67	0.257	0.0770	1.70	3.10
0.7	80	4.00	1.02	11.08	0.271	0.0920	1.70	3.06
0.7	90	4.38	1.09	11.08	0.271	0.0980	1.70	2.99
0.7	100	4.81	1.16	11.08	0.271	0.1050	1.70	2.89

TABLE 2-Continued

f(Hz)	Span(%)	e(in.)	H(ft)	L(ft)	d/L	H/L	$2\pi d/L$	H/e
0.8	10	0.41	0.14	8.33	0.360	0.017	2.18	4.67
0.8	20	0.88	0.30	8.33	0.360	0.036	2.18	4.29
0.8	30	1.31	0.42	8.67	0.346	0.048	2.18	3.82
0.8	40	1.81	0.55	8.75	0.343	0.063	2.18	3.67
0.8	50	2.25	0.64	9.10	0.330	0.070	2.18	3.37
0.8	60	2.69	0.74	8.75	0.343	0.085	2.18	3.36
0.9	10	0.37	0.13	6.53	0.4594	0.020	2.81	4.33
0.9	20	0.75	0.28	6.28	0.4778	0.045	2.81	4.67
0.9	30	1.22	0.42	6.53	0.4594	0.064	2.81	4.20
0.9	40	1.63	0.50	7.32	0.4098	0.068	2.81	3.57
0.9	50	2.06	0.64	7.32	0.4098	0.087	2.81	3.76
1.0	10	0.31	0.13	6.36	0.472	0.020	3.45	4.33
1.0	20	0.69	0.25	5.47	0.548	0.046	3.45	4.17
1.0	30	1.06	0.38	5.47	0.548	0.069	3.45	4.22
1.0	40	1.47	0.48	5.60	0.536	0.086	3.45	4.00

TABLE 3

Attenuation Coefficients

$f(\text{Hz})$	$d(\text{ft})$	$H_1(\text{ft})$	$H_2(\text{ft})$	$L(\text{ft})$	$K_{\text{mea.}} (\text{ft}^{-1}) = \ln(H_1/H_2)/61.8\text{ft}$	$K_{\text{theo.}} (\text{ft}^{-1})$
0.4	3	1.031	1.0267	22.070	$6.763 \times 10^{-5}$	$8.487 \times 10^{-5}$
0.4	3	1.187	1.1846	23.804	$3.275 \times 10^{-5}$	$7.973 \times 10^{-5}$
0.4	3	0.942	0.9390	23.804	$4.302 \times 10^{-5}$	$7.973 \times 10^{-5}$
0.5	3	0.731	0.7298	18.025	$3.544 \times 10^{-5}$	$8.889 \times 10^{-5}$
0.5	3	0.850	0.8470	17.050	$5.720 \times 10^{-5}$	$9.258 \times 10^{-5}$
0.5	3	0.990	0.9854	18.847	$7.536 \times 10^{-5}$	$8.596 \times 10^{-5}$
0.6	3	0.798	0.7940	13.510	$8.131 \times 10^{-5}$	$9.882 \times 10^{-5}$
0.6	3	0.924	0.9190	13.730	$8.780 \times 10^{-5}$	$9.782 \times 10^{-5}$

TABLE 4  
Model Calibration

Location of Load(ft)	Load(lb)									
	2	4	6	8	10	12	14	16	18	20
2.500	12	24	36	49	62	74	86	99	110	123
2.333	12	23	35	46	57	68	80	91	102	114
2.167	10	21	31	42	52	63	73	84	95	105
2.000	10	20	30	39	49	59	68	78	88	98
1.917	10	19	29	38	48	57	66	76	85	94
1.833	9	18	27	36	45	54	63	72	81	90
1.750	8	17	26	34	43	52	60	69	77	86
1.667	8	16	24	33	41	49	58	66	73	82
1.583	8	16	24	32	39	47	55	63	70	78
1.500	7	15	22	30	37	45	52	59	67	74
1.417	7	14	21	28	35	42	49	55	63	70
1.333	7	13	20	27	33	40	47	53	60	66
1.250	6	12	18	24	30	36	42	49	55	61
1.000	4	9	14	19	24	29	34	39	44	49

Inset values are strains ( $\times 10^{-6}$ ).

Location of load is based on a datum level of 0 ft at the wave tank floor.

**TABLE 5**

**Measured Wave Forces and Moments**





Run	$B_{hm}$	$B_{hf}$	$M_{S_1}$ (lb-ft)	$F_h$ (lb)	$MF_h$ (lb)	$L_1$ (ft)
1	0.000	25.000	0.706	0.478	0.571	1.478
2	0.000	9.000	1.713	1.131	1.429	1.515
3	0.000	9.000	3.284	2.119	2.500	1.530
4	0.000	9.000	4.767	2.989	3.750	1.595
5	0.000	9.000	6.194	3.835	5.250	1.615
6	0.000	9.000	8.009	4.945	6.753	1.653
7	0.000	9.000	9.394	5.655	8.485	1.691
8	38.365	41.199	12.201	6.060	10.000	1.753
9	35.496	37.674	14.931	8.392	11.547	1.779
10	29.002	30.955	16.409	9.026	12.679	1.818
11	90.000	90.000	11.059	6.689	9.750	1.537
12	57.336	66.343	2.864	1.793	2.000	1.593
13	90.000	90.000	4.132	2.533	3.250	1.627
14	90.000	90.000	6.971	4.253	5.000	1.639
15	69.017	90.000	8.124	4.735	6.480	1.716
16	64.445	81.242	9.724	5.528	7.765	1.759
17	47.422	53.826	11.083	6.067	8.220	1.827
18	90.000	90.000	14.606	8.086	10.000	1.806
19	72.658	90.000	14.848	7.995	11.111	1.857
20	41.810	47.435	17.445	8.799	12.200	1.983
21	90.000	90.000	11.050	6.653	9.750	1.622
22	89.355	90.000	3.178	1.908	2.250	1.666
23	90.000	90.000	6.003	3.496	4.000	1.717
24	90.000	90.000	9.132	5.203	6.000	1.765
25	90.000	90.000	10.947	5.969	7.333	1.834
26	90.000	90.000	12.713	6.702	8.216	1.872
27	72.247	90.000	14.143	7.442	9.792	1.900
28	90.000	90.000	17.657	8.838	11.166	1.998
29	35.100	41.311	18.796	8.764	11.430	2.145
30	35.120	46.498	20.737	9.635	12.570	2.152
31	90.000	90.000	11.413	5.832	7.000	1.695
32	90.000	90.000	4.096	2.345	2.823	1.747
33	90.000	90.000	7.946	4.415	5.111	1.800
34	90.000	90.000	10.594	5.706	6.667	1.857
35	90.000	90.000	12.219	6.410	7.579	1.906
36	90.000	90.000	14.479	7.514	8.750	1.927
37	90.000	90.000	14.302	7.750	9.375	1.946
38	90.000	90.000	17.304	8.574	11.186	2.018
39	90.000	90.000	17.657	8.633	12.203	2.045
40	90.000	90.000	20.482	9.933	12.368	2.072

Run.	$M_{ST}$ (1b-ft)	C <sub>L</sub>	F <sub>T</sub> (1b)	M <sub>T</sub> (1b)	$M_{ST}$ (1b-ft), L2(ft)
1	0.000	0.000	0.000	0.000	0.000
2	0.000	0.000	0.000	0.000	0.000
3	0.000	0.000	0.000	0.000	0.000
4	0.000	0.000	0.000	0.000	0.000
5	0.000	0.000	0.000	0.000	0.000
6	0.353	0.036	0.073	0.215	0.127
7	1.413	0.101	0.283	0.886	0.509
8	1.413	0.119	0.452	0.886	0.819
9	1.766	0.120	0.597	1.111	1.095
10	2.625	0.170	1.002	1.768	1.865
11	0.000	0.000	0.000	0.000	0.000
12	0.000	0.000	0.000	0.000	0.000
13	0.000	0.000	0.000	0.000	0.000
14	0.000	0.000	0.000	0.000	0.000
15	0.353	0.042	0.087	0.200	0.162
16	0.353	0.030	0.085	0.200	0.162
17	0.706	0.050	0.199	0.361	0.338
18	1.059	0.043	0.214	0.540	0.424
19	1.766	0.059	0.380	0.905	0.777
20	2.119	0.060	0.570	1.091	1.208
21	0.000	0.000	0.000	0.000	0.000
22	0.000	0.000	0.000	0.000	0.000
23	0.000	0.000	0.000	0.000	0.000
24	0.000	0.000	0.000	0.000	0.000
25	0.000	0.000	0.000	0.000	0.000
26	0.353	0.021	0.063	0.178	0.134
27	0.706	0.034	0.144	0.358	0.311
28	1.413	0.037	0.259	0.615	0.593
29	1.766	0.062	0.504	0.769	1.165
30	2.119	0.060	0.568	0.923	1.335
31	0.000	0.000	0.000	0.000	0.000
32	0.000	0.000	0.000	0.000	0.000
33	0.000	0.000	0.000	0.000	0.000
34	0.000	0.000	0.000	0.000	0.000
35	0.353	0.010	0.022	0.167	0.049
36	0.353	0.009	0.028	0.167	0.064
37	0.706	0.022	0.091	0.333	0.205
38	1.059	0.027	0.139	0.461	0.328
39	1.059	0.023	0.142	0.461	0.339
40	1.766	0.042	0.306	0.769	0.742

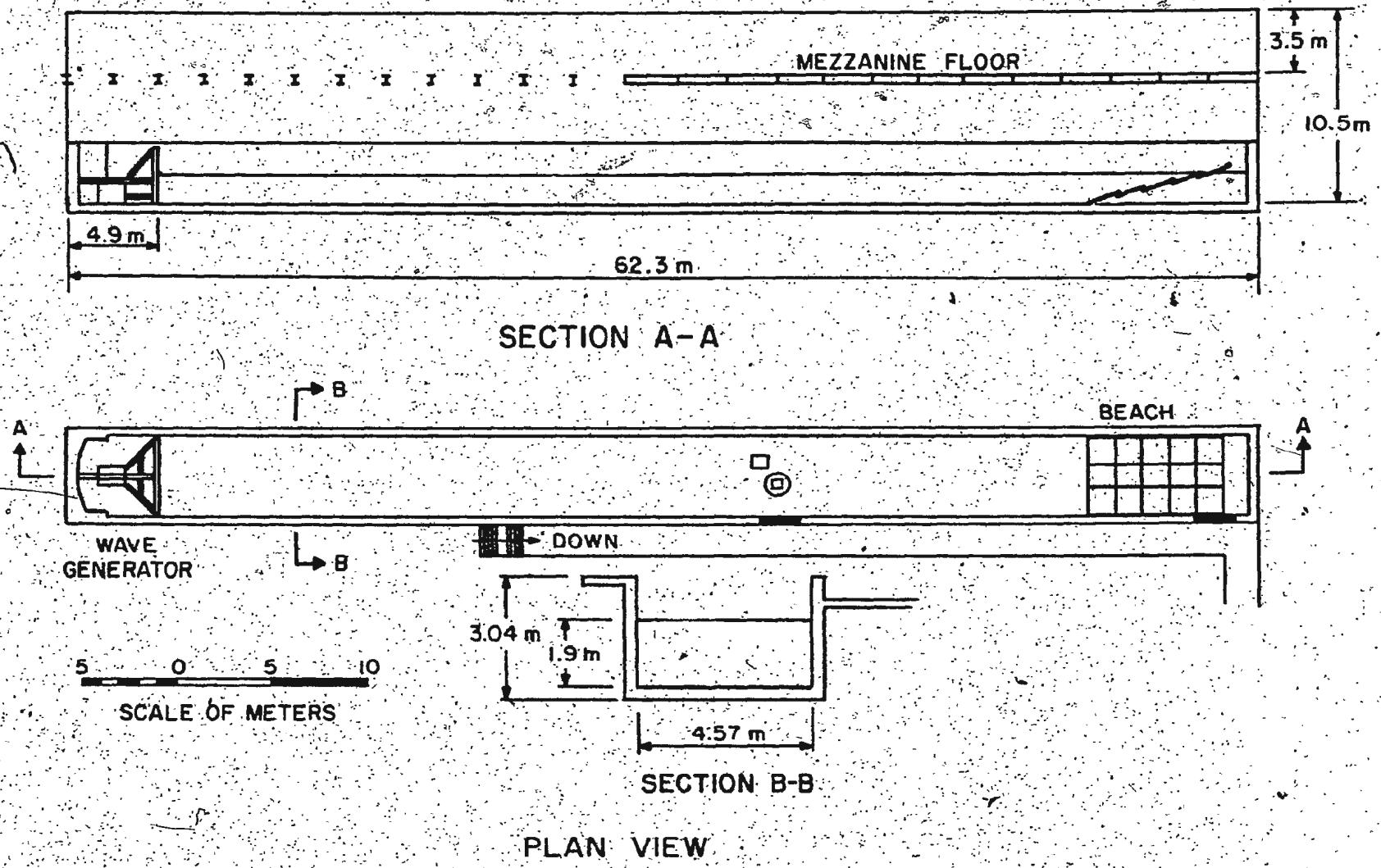


Fig. I: Elevation and Plan Views of Wave Tank

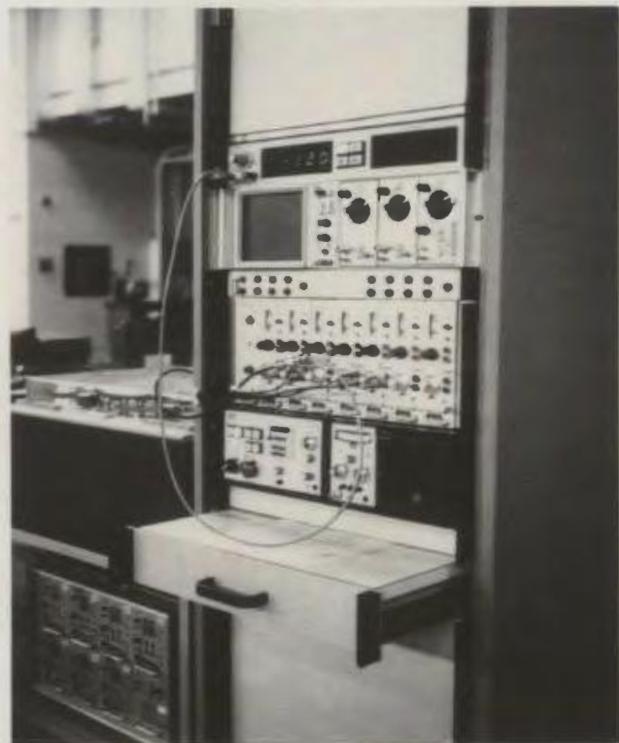


Fig. 2: Wave Generator Control Console

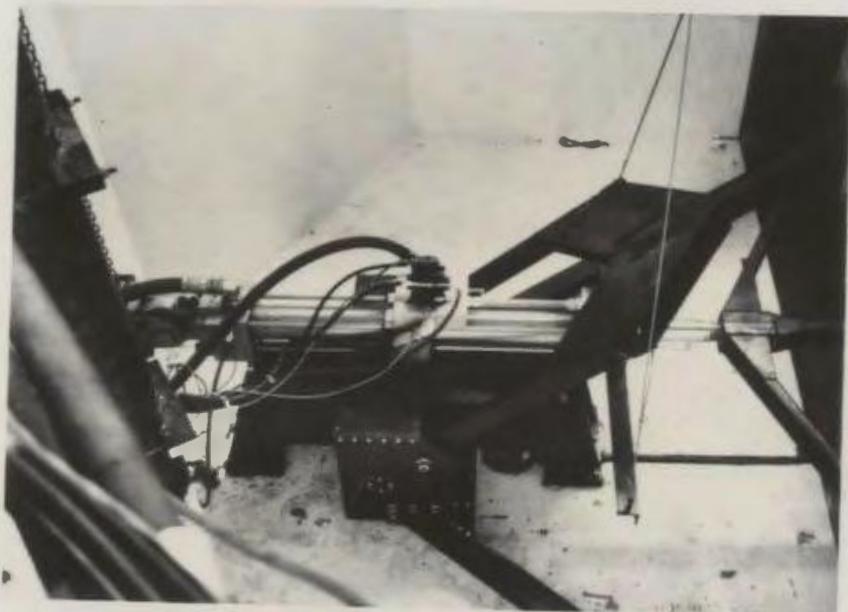


Fig. 3: Hydraulic Actuator

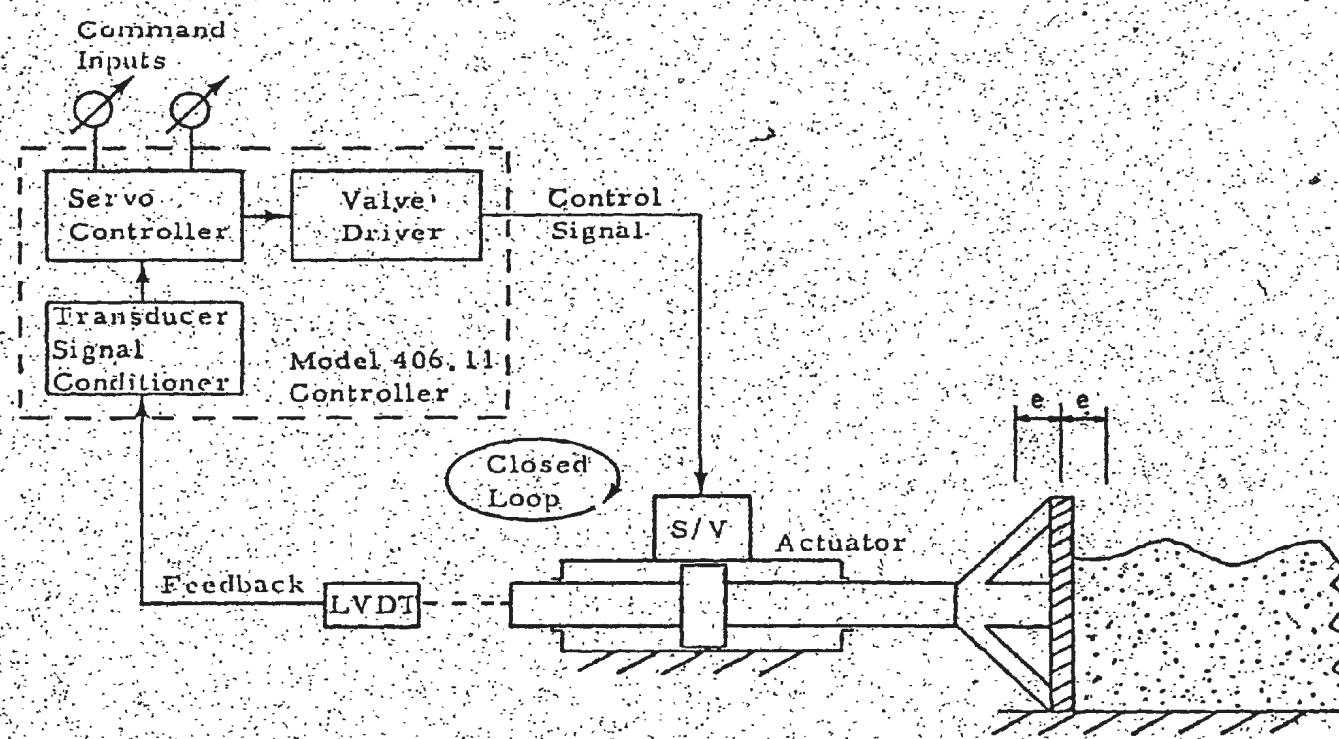


Fig. 4: Control Loop for Piston-Type Wave Generator



Fig. 5: Wave Absorbing Beach (Slope 1:6.2)



Fig. 6: General View of Wave Tank

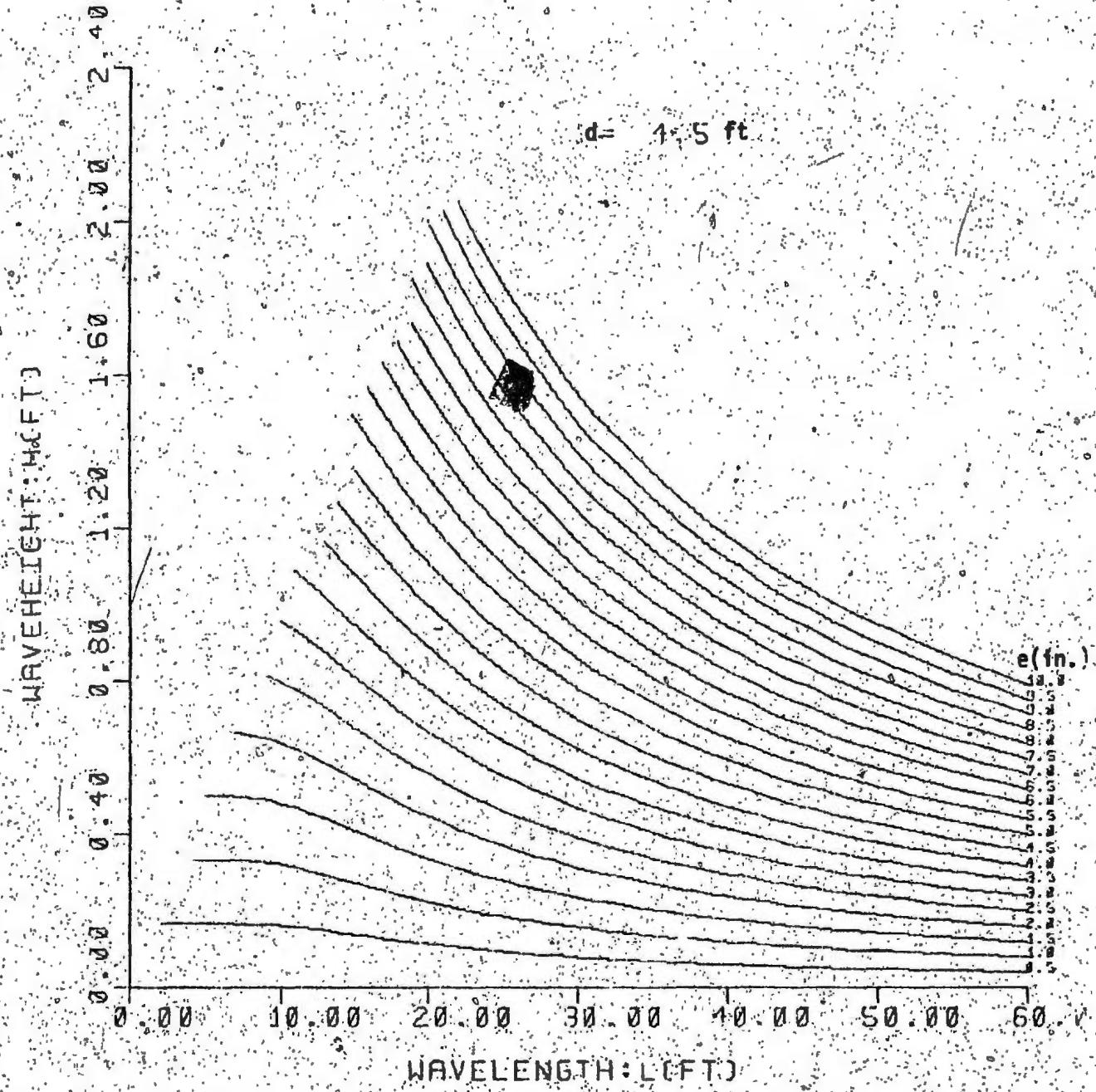


Fig. 7: Theoretical Performance Envelope for 4.5 ft Water Depth

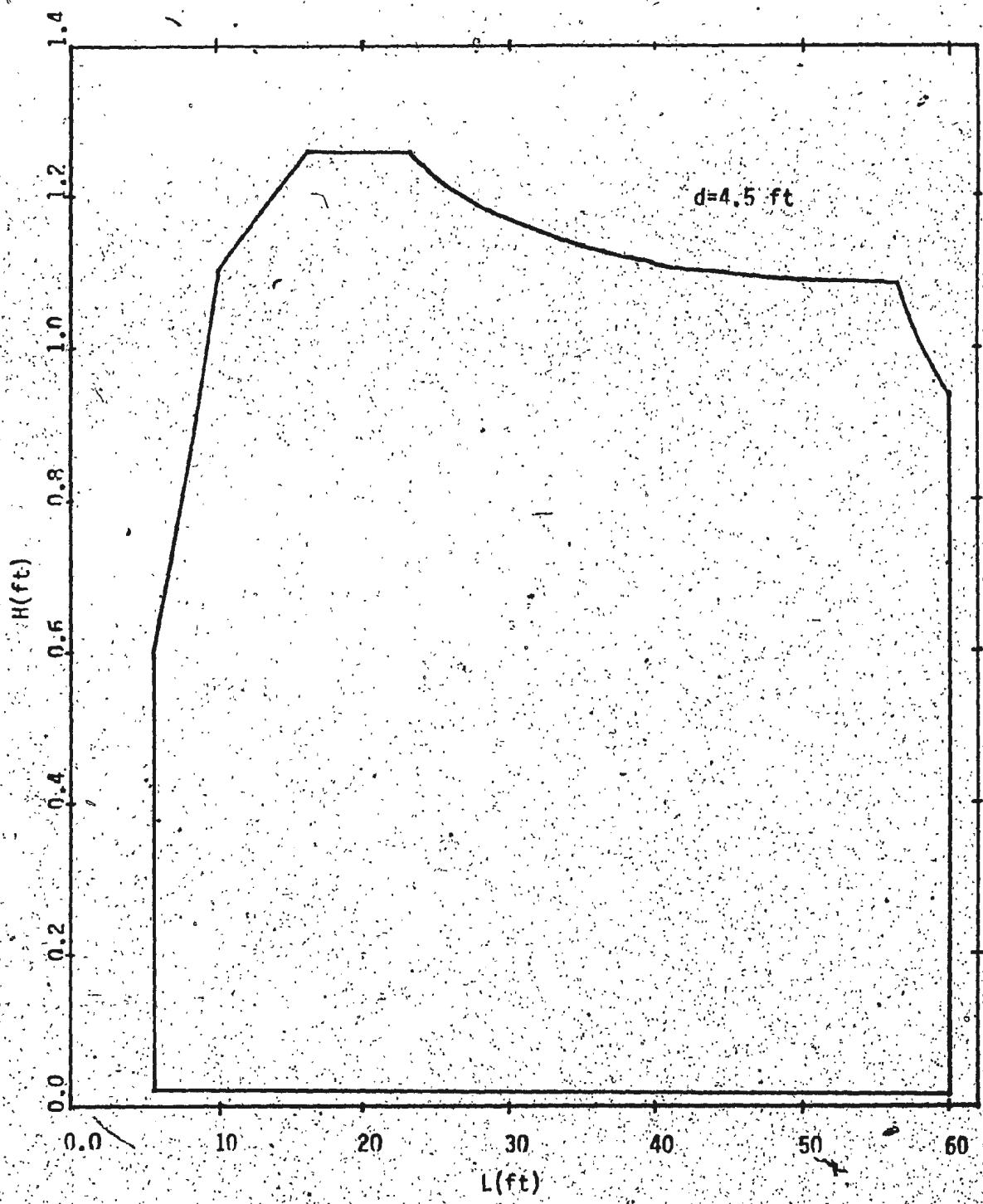


Fig. 8: Experimental Performance Envelope for 4.5 ft Water Depth

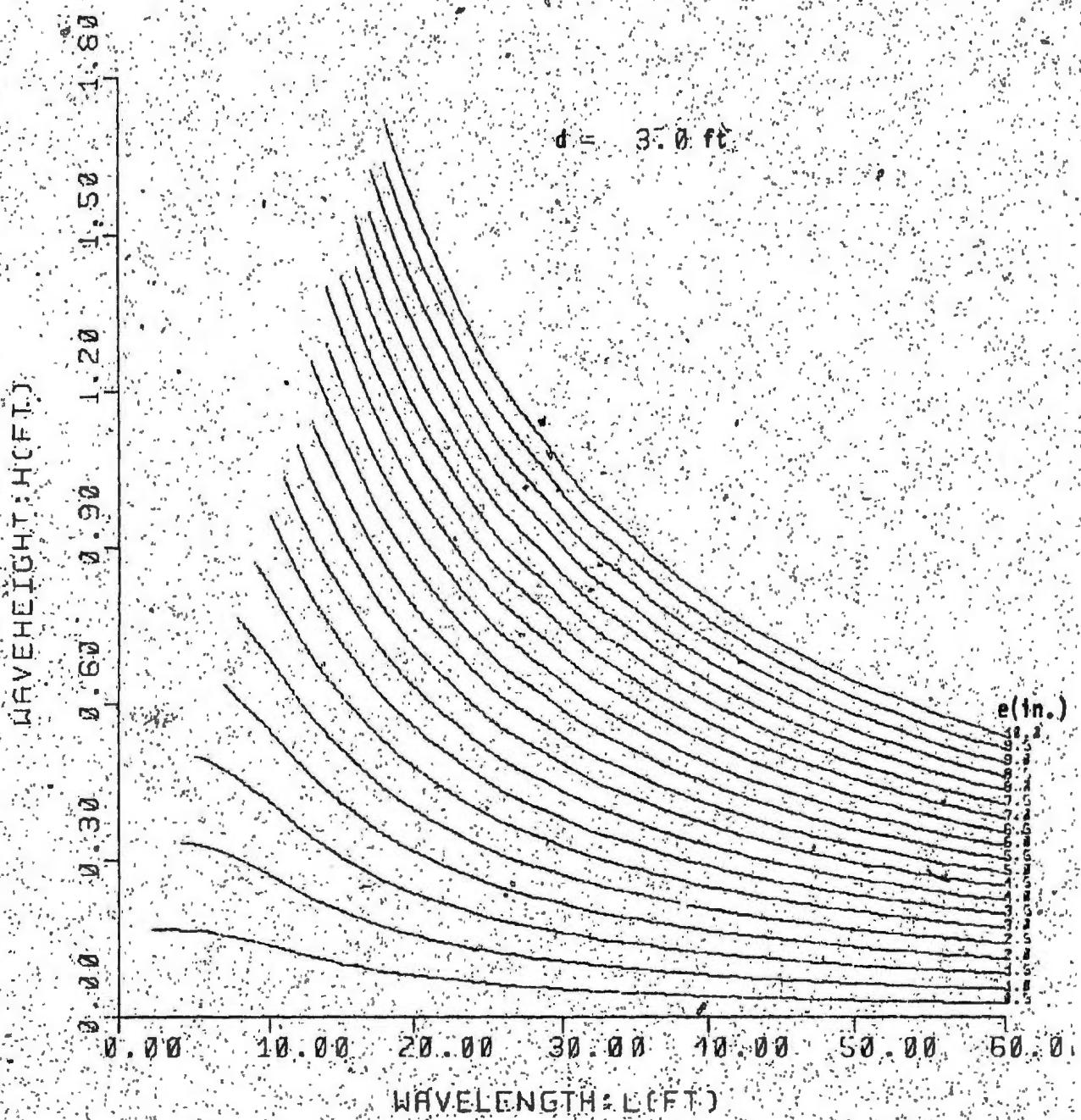


Fig. 9: Theoretical Performance Envelope for 3 ft Water Depth

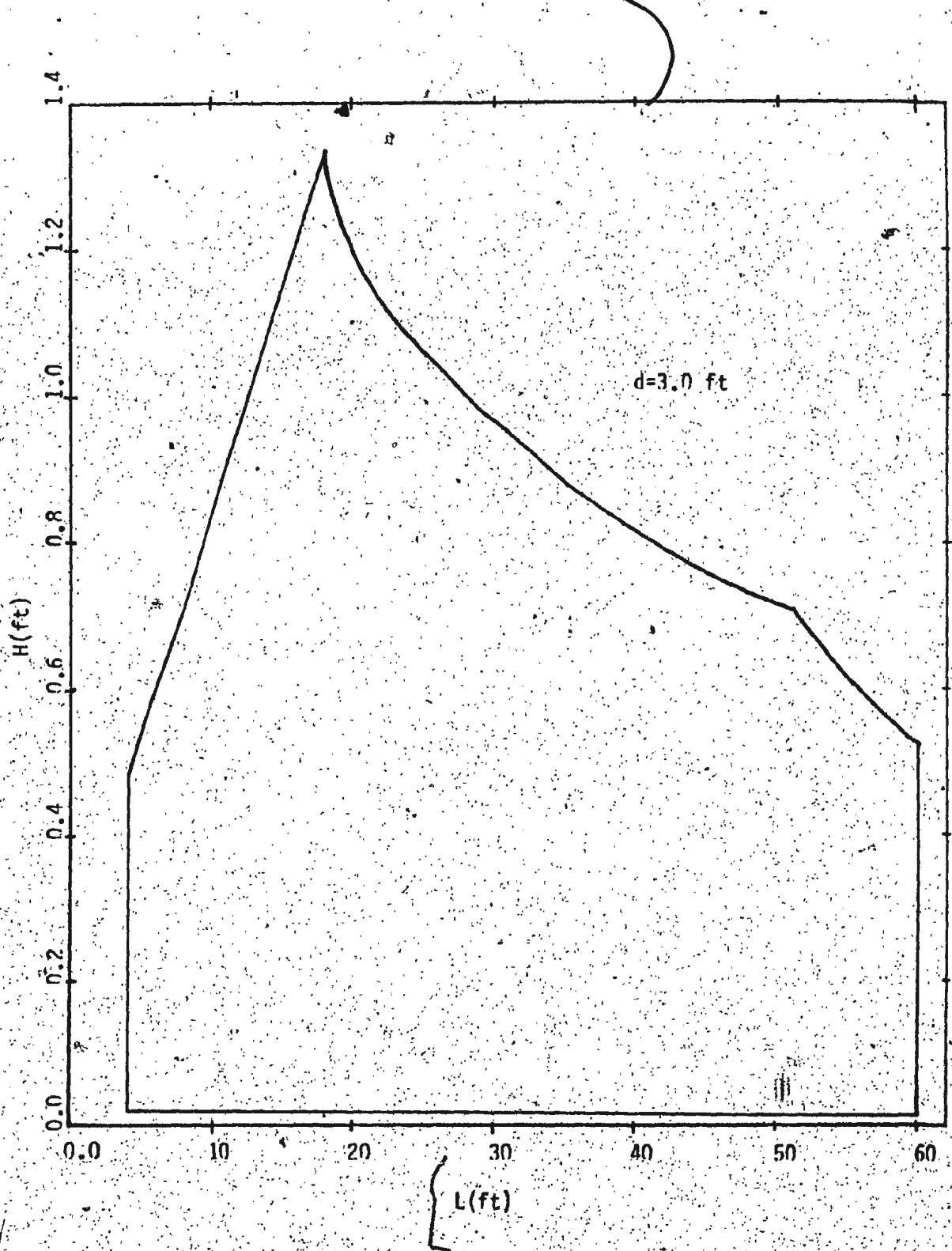


Fig. 10: Experimental Performance Envelope for 3 ft Water Depth

**Experiment**

**Theory**

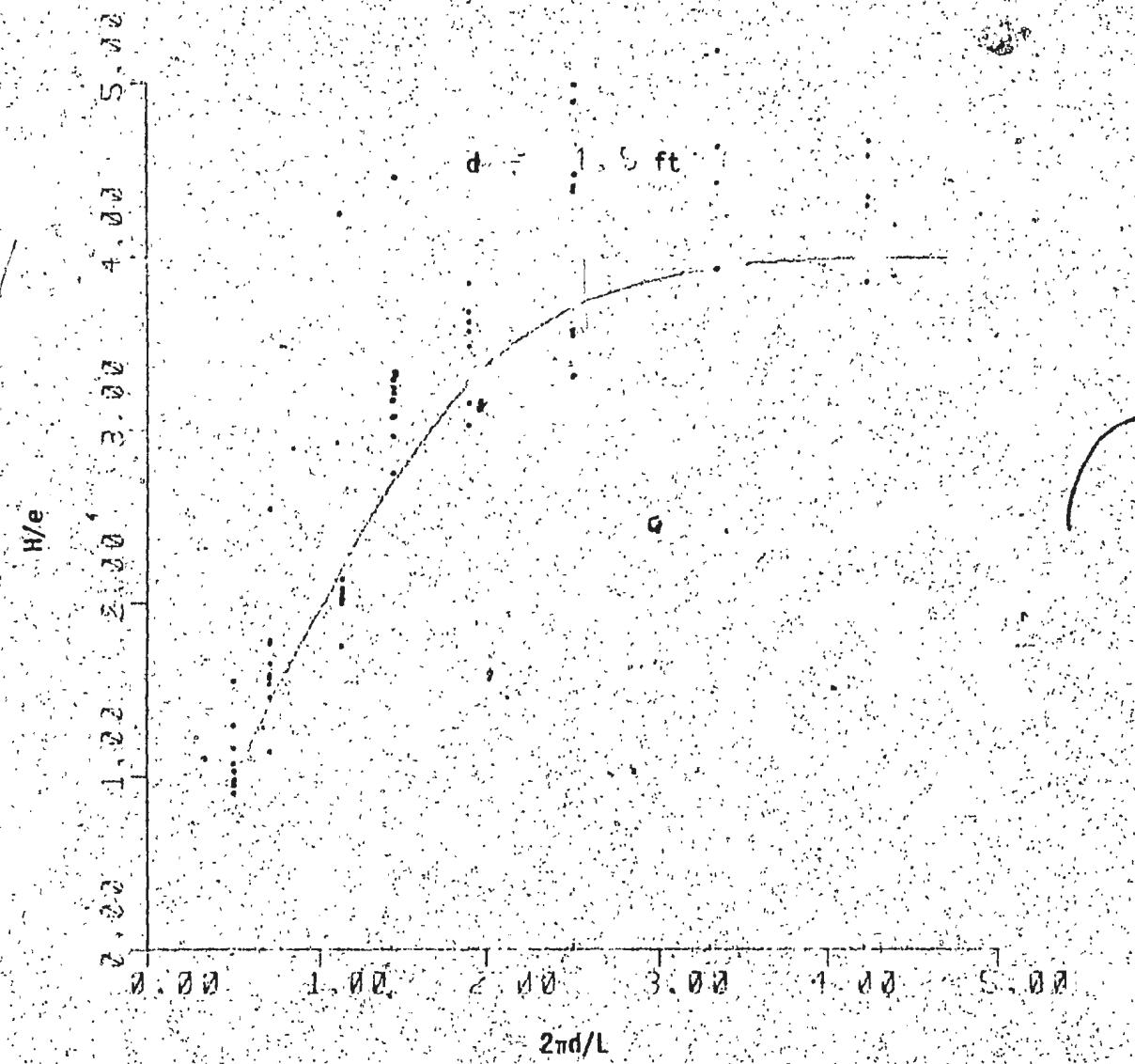


Fig. 11: Wave Generator Transfer Function for 4.5 ft Water Depth

Experiment

Theory

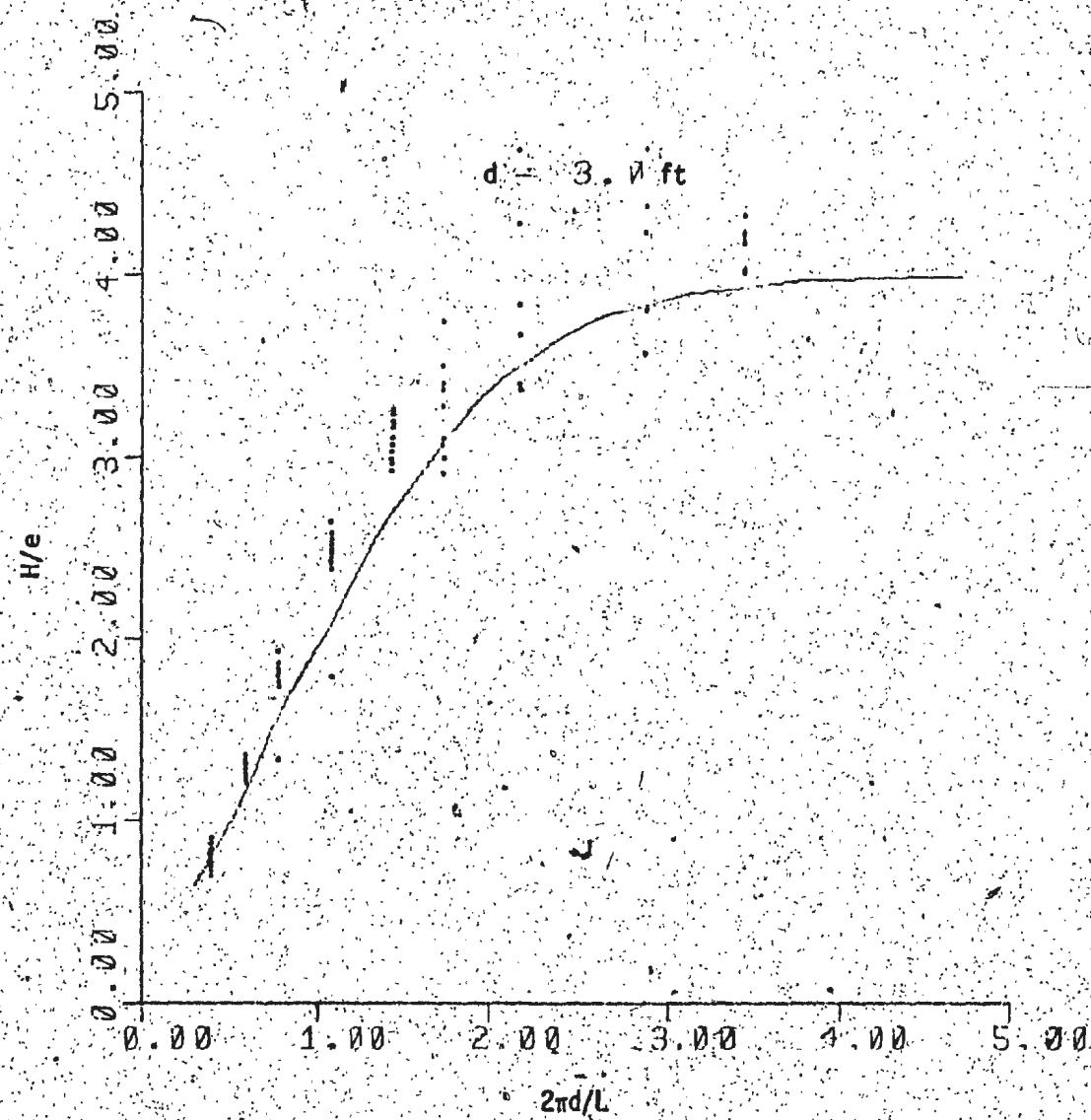


Fig. 12: Wave Generator Transfer Function for 3 ft Water Depth

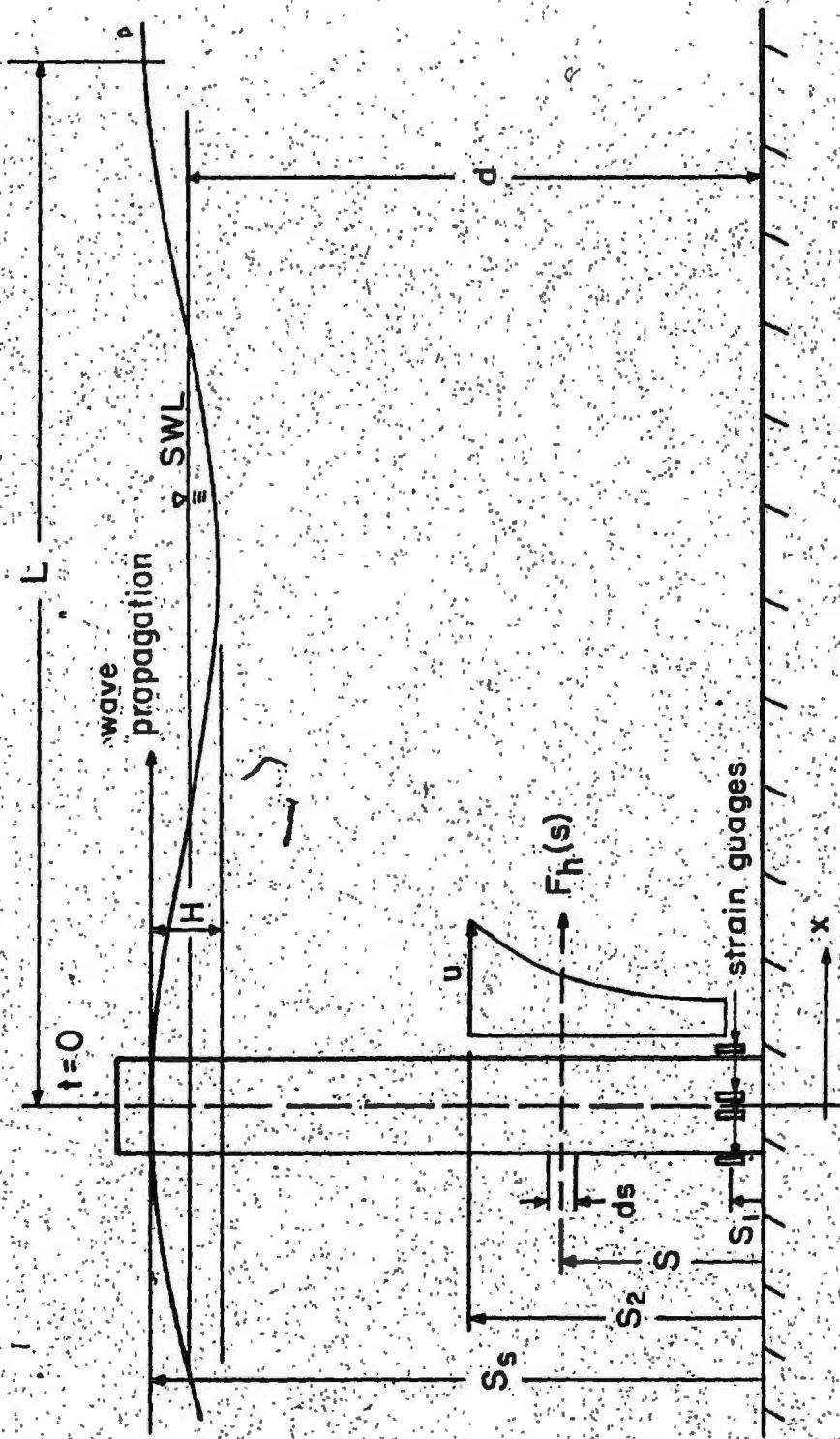


Fig. 13: Definition Sketch



Fig. 14: View of Wave Gauges Used to Measure Reflection Coefficients

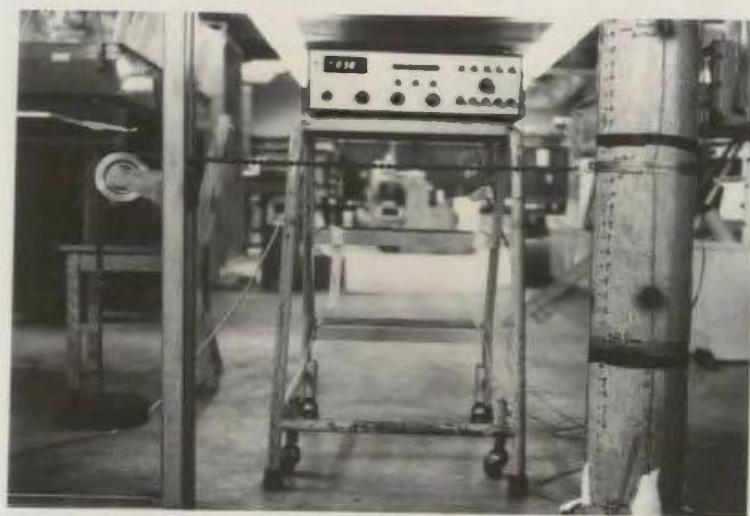


Fig. 15: Calibration of Model

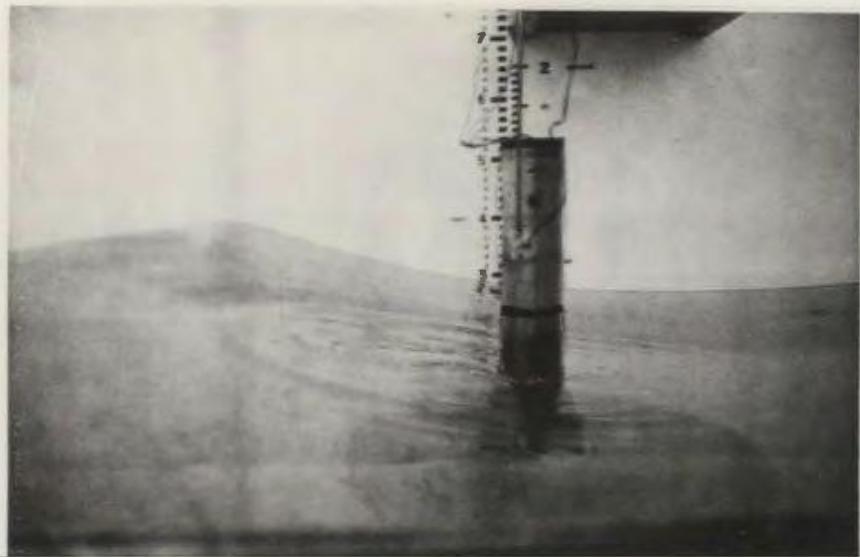


Fig. 16: Cylinder Under Test



Fig. 17: Vortex Shedding in the Wake of the Circular Cylinder

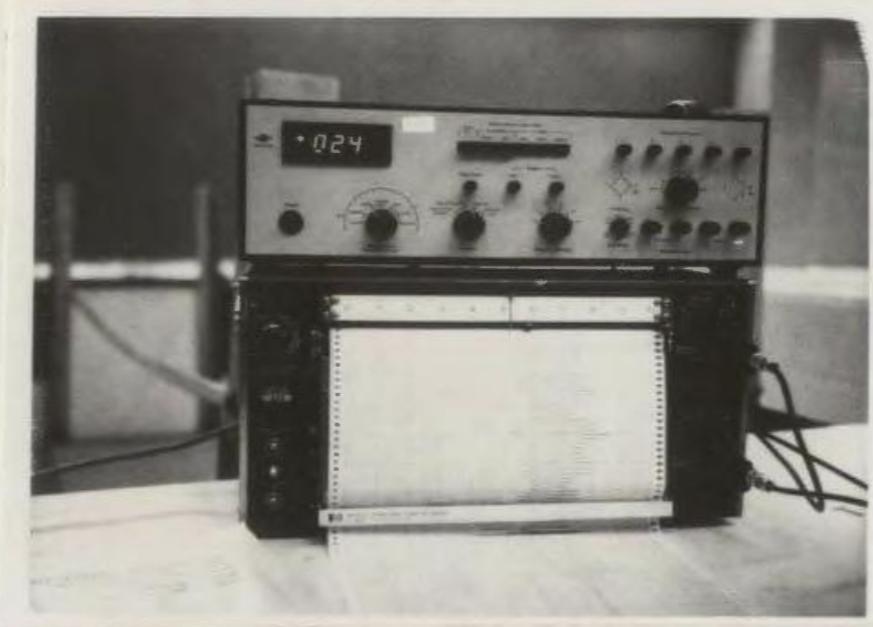


Fig. 18: View of Typical Output from the Strain Gage Bridges



Fig. 19: Tensile Specimen in Testing Machine

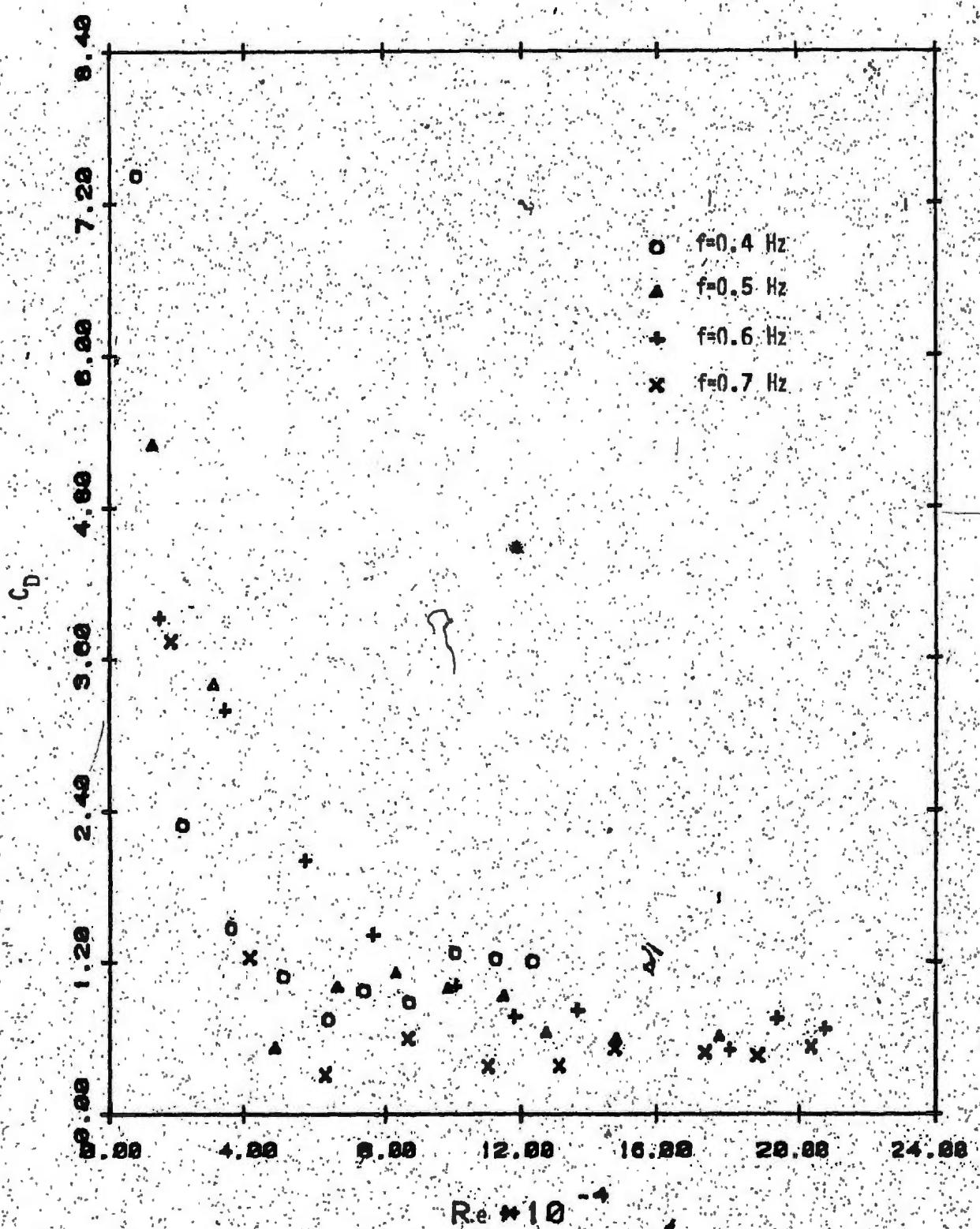


Fig. 20: Drag coefficient vs. Reynolds Number

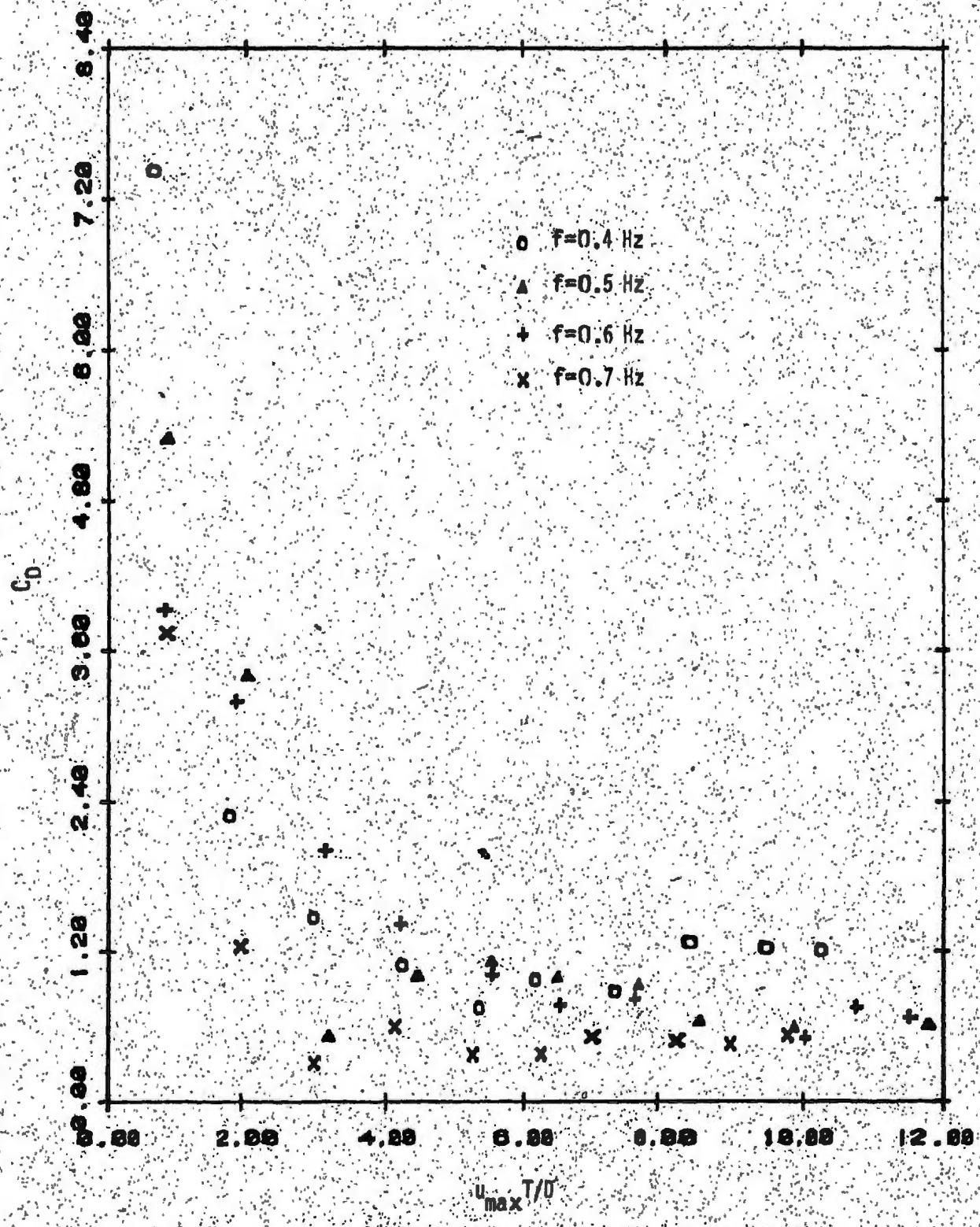


Fig. 21: Drag Coefficient vs. Keulegan-Carpenter Number

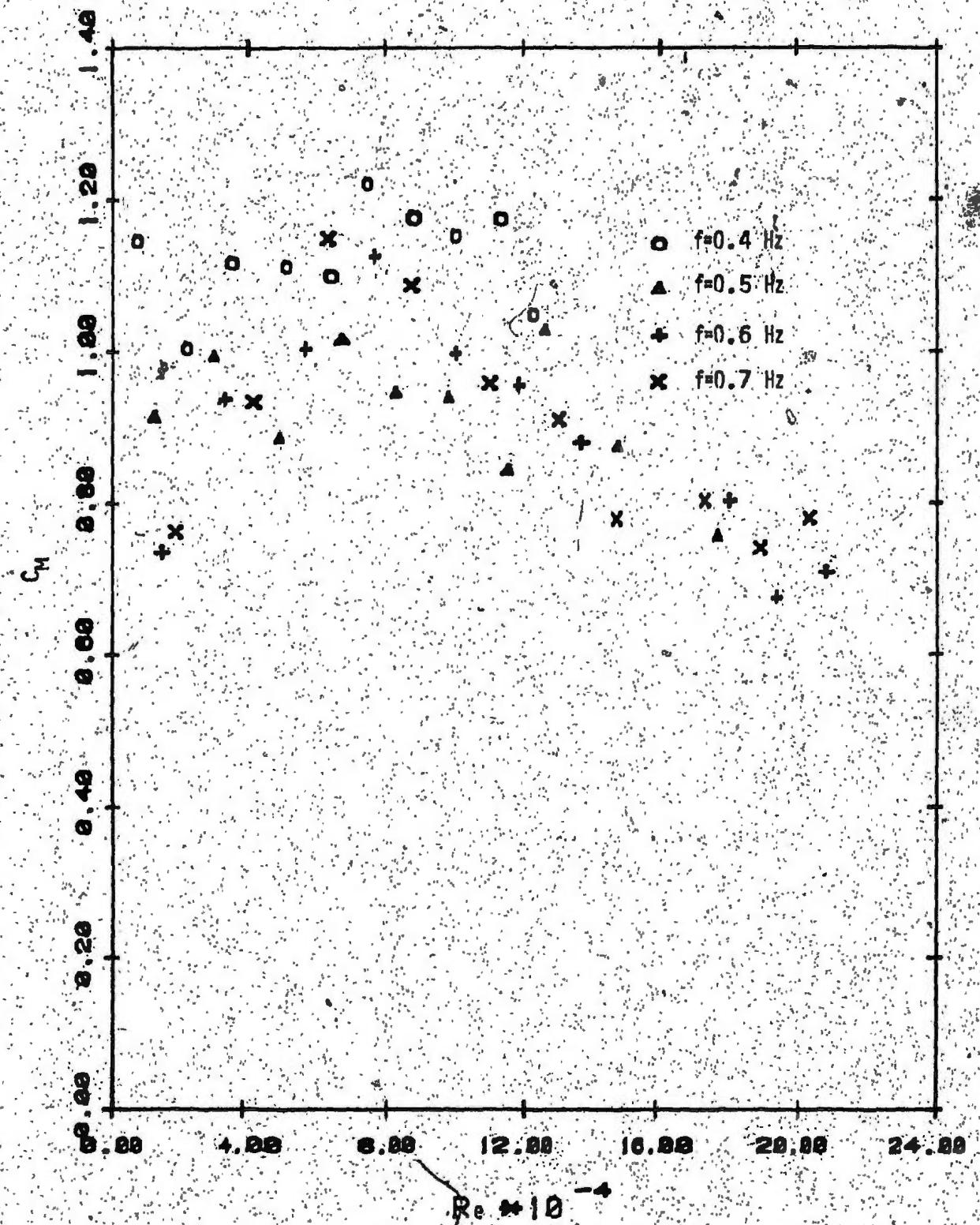


Fig. 22: Mass Coefficient vs. Reynolds Number

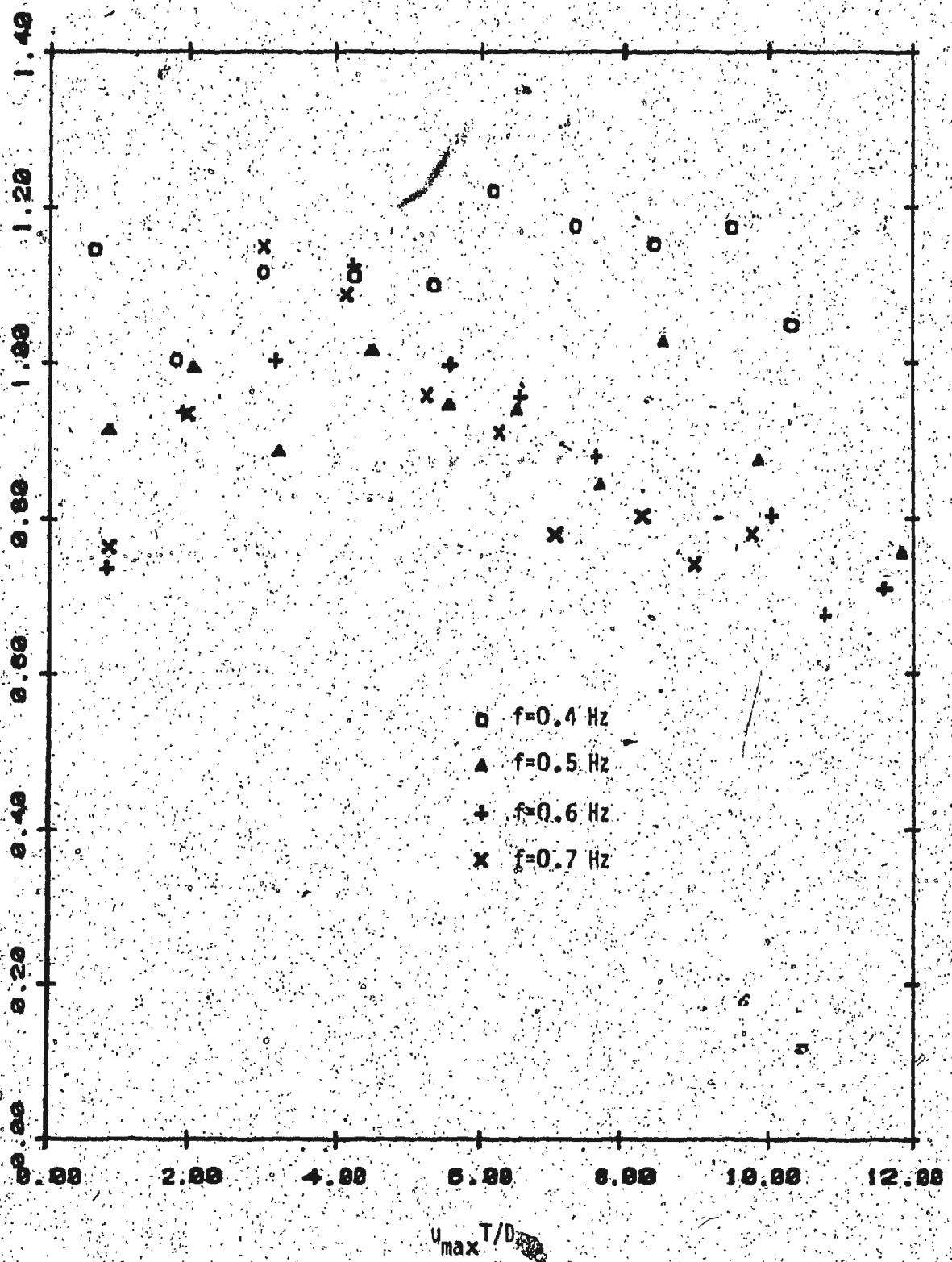


Fig. 23: Mass Coefficient vs. Keulegan-Carpenter Number

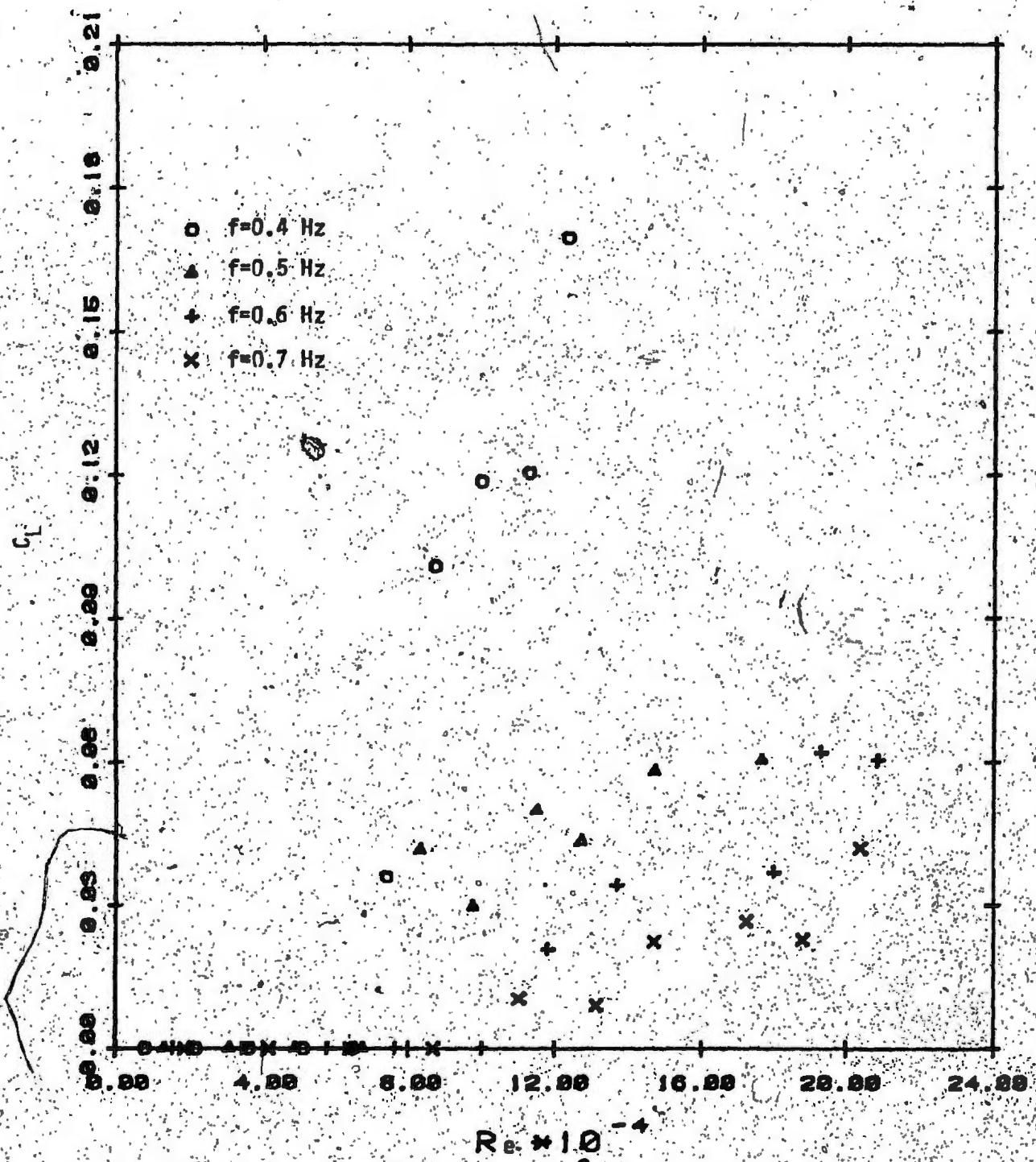


Fig. 24: Lift Coefficient-vs. Reynolds Number

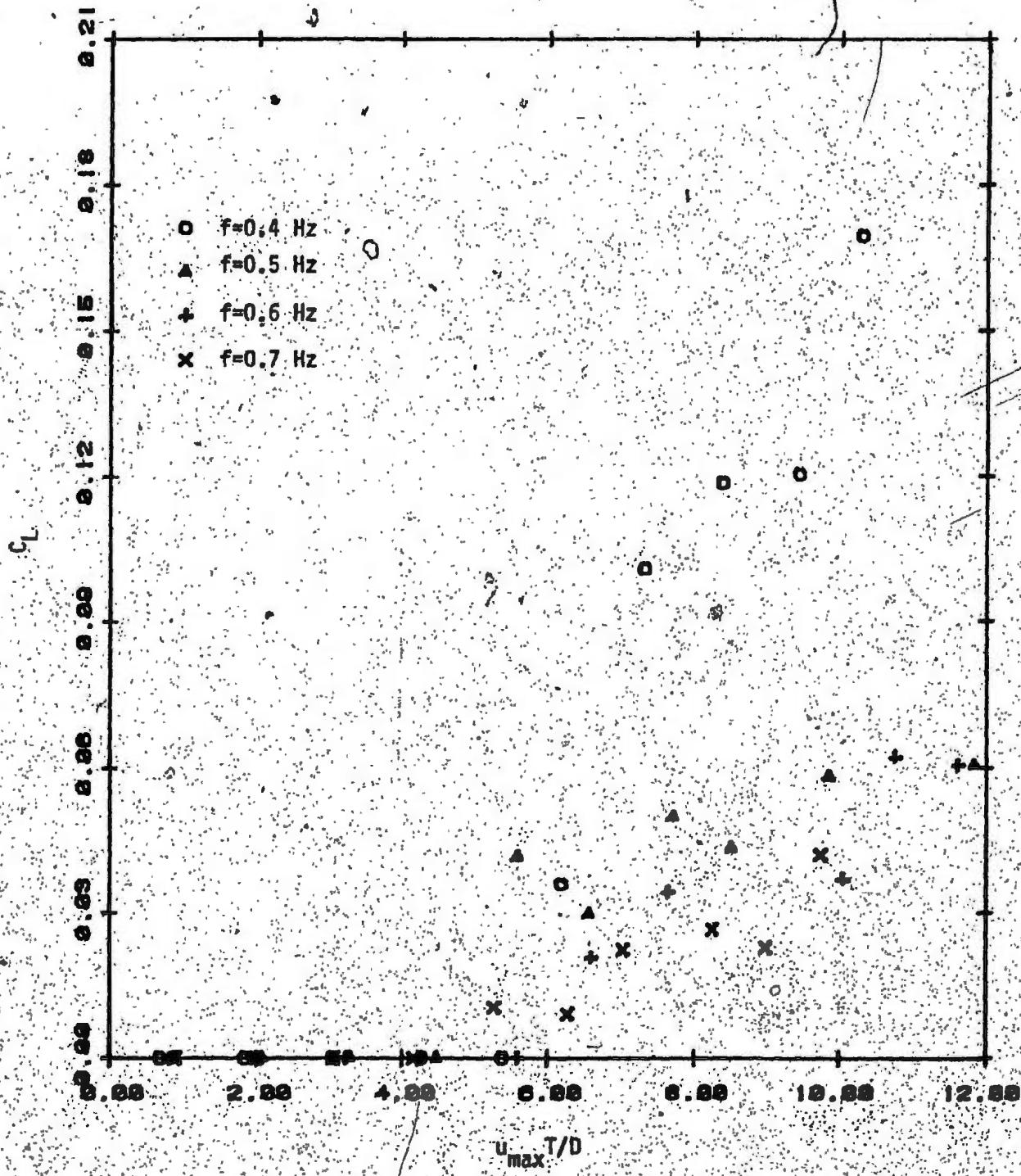


Fig. 25: Lift Coefficient vs. Keulegan-Carpenter Number

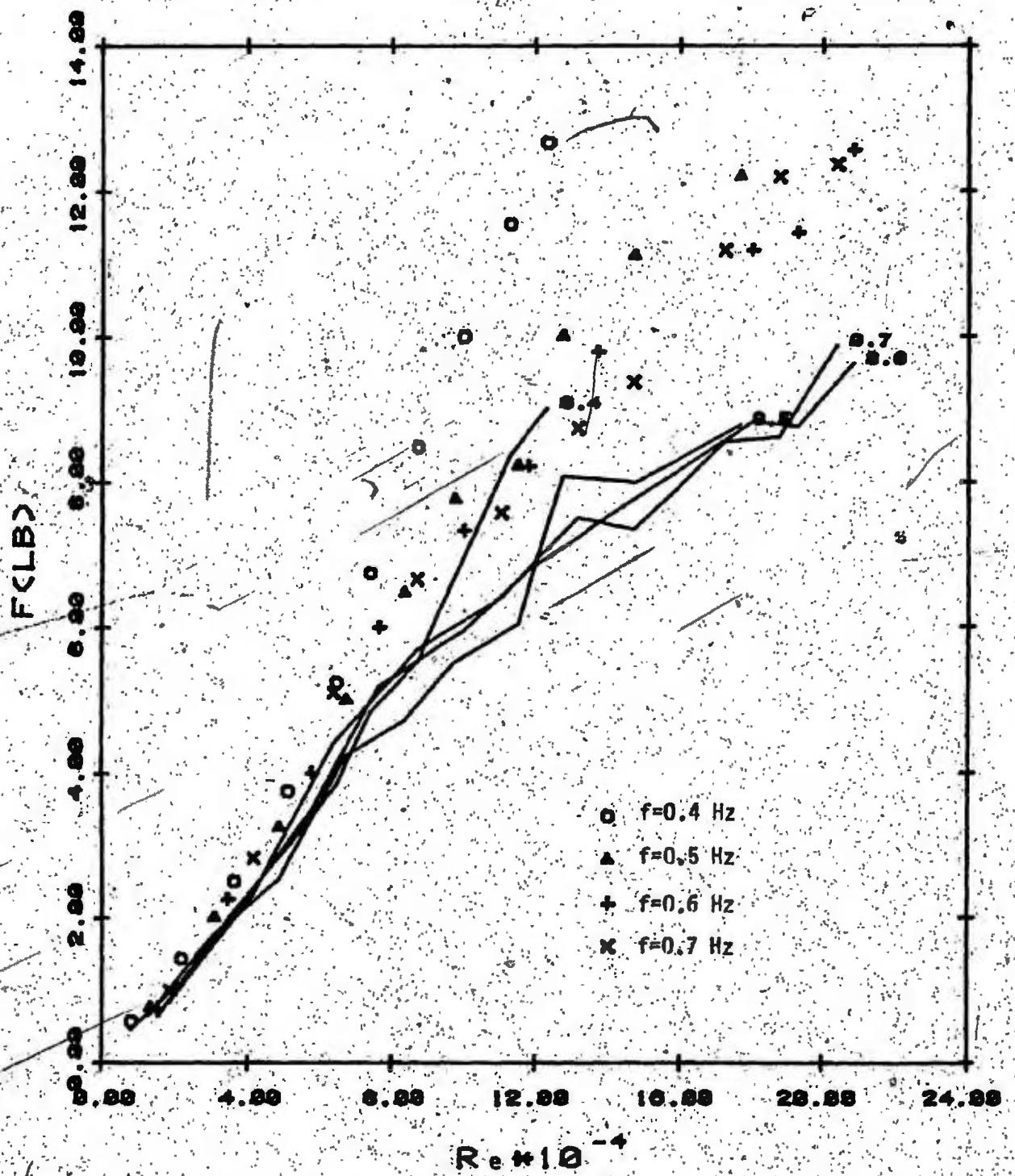


Fig. 26: Longitudinal Force vs. Reynolds Number

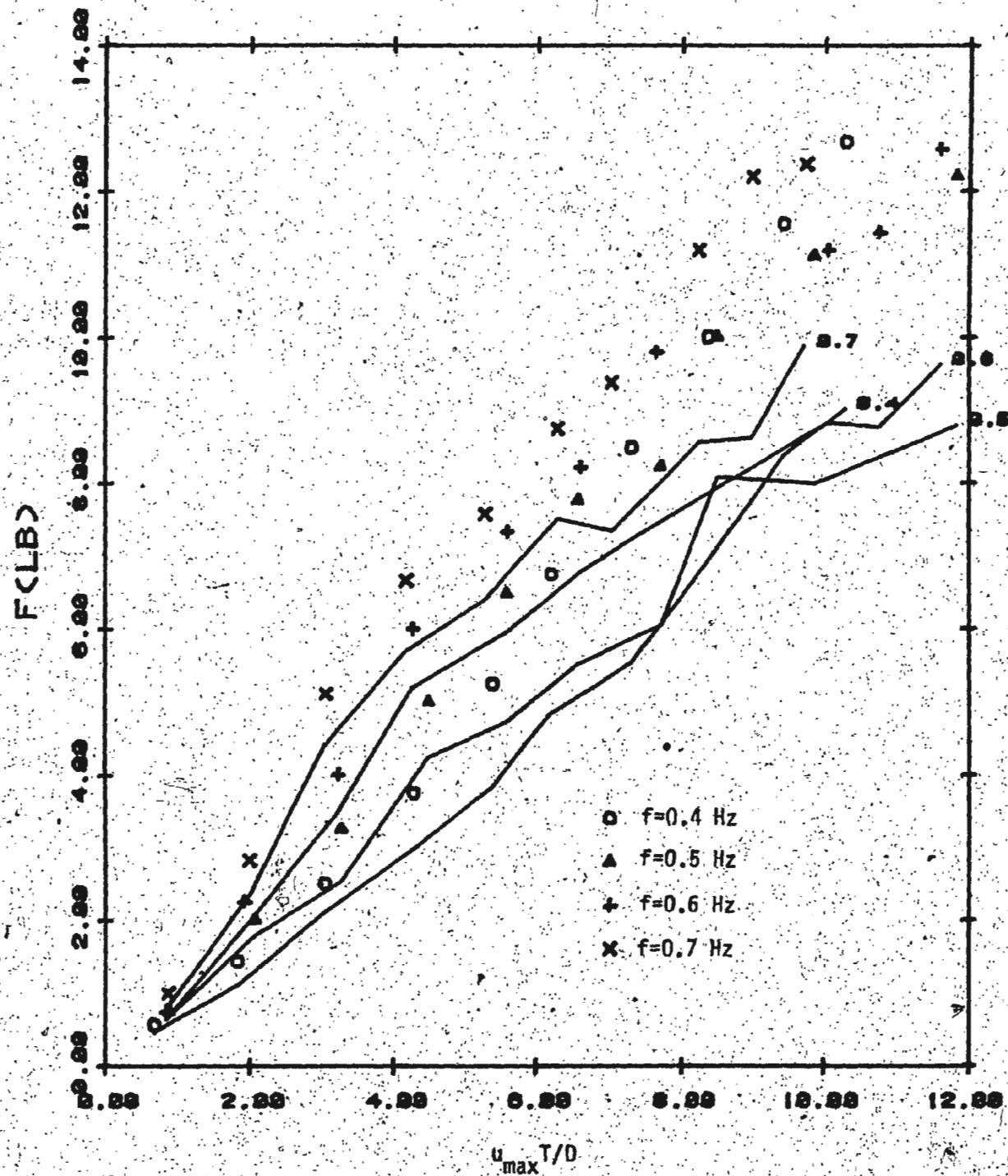


Fig. 27: Longitudinal Force vs. Keulegan-Carpenter Number

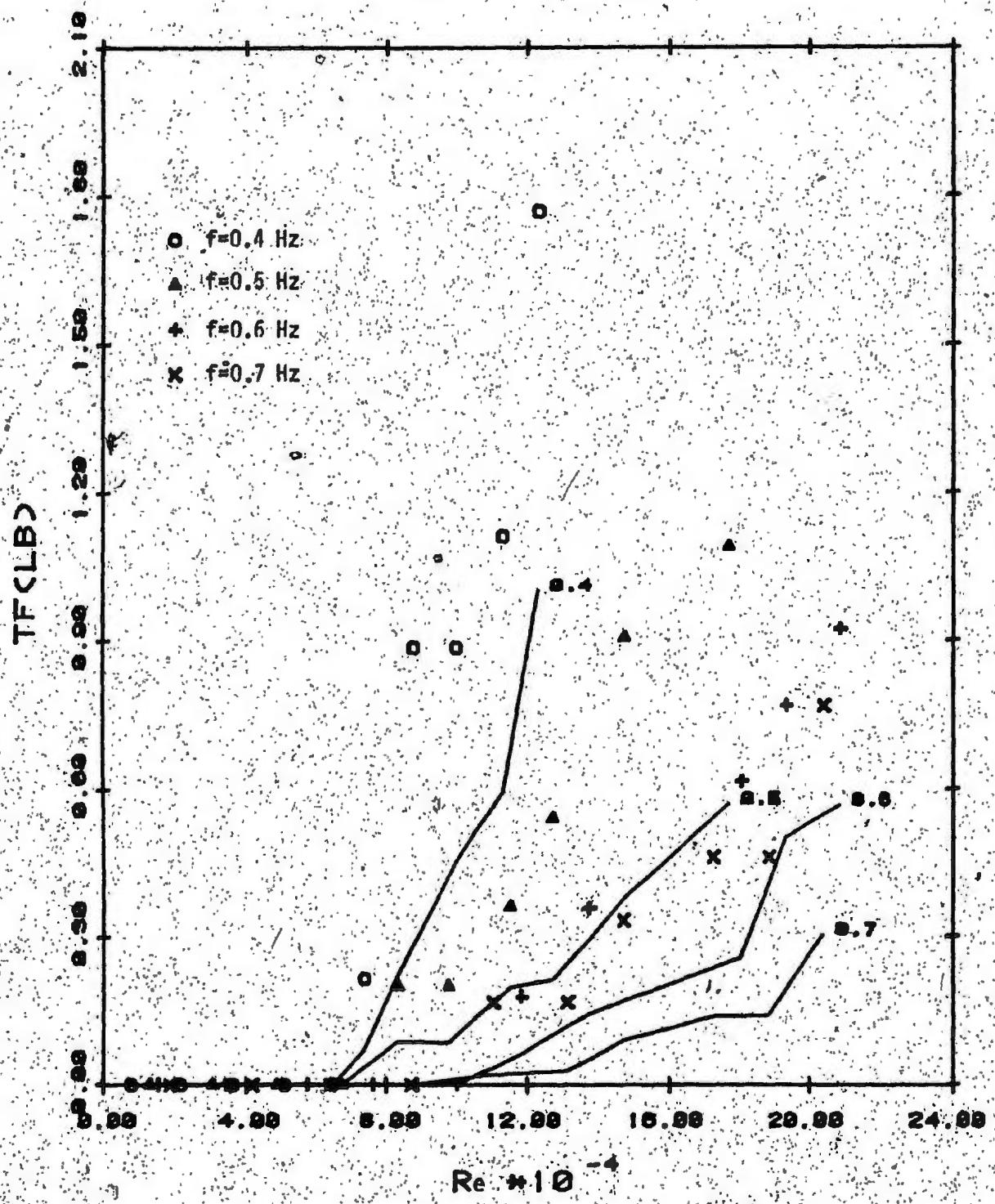


Fig. 28: Transverse Force vs. Reynolds Number

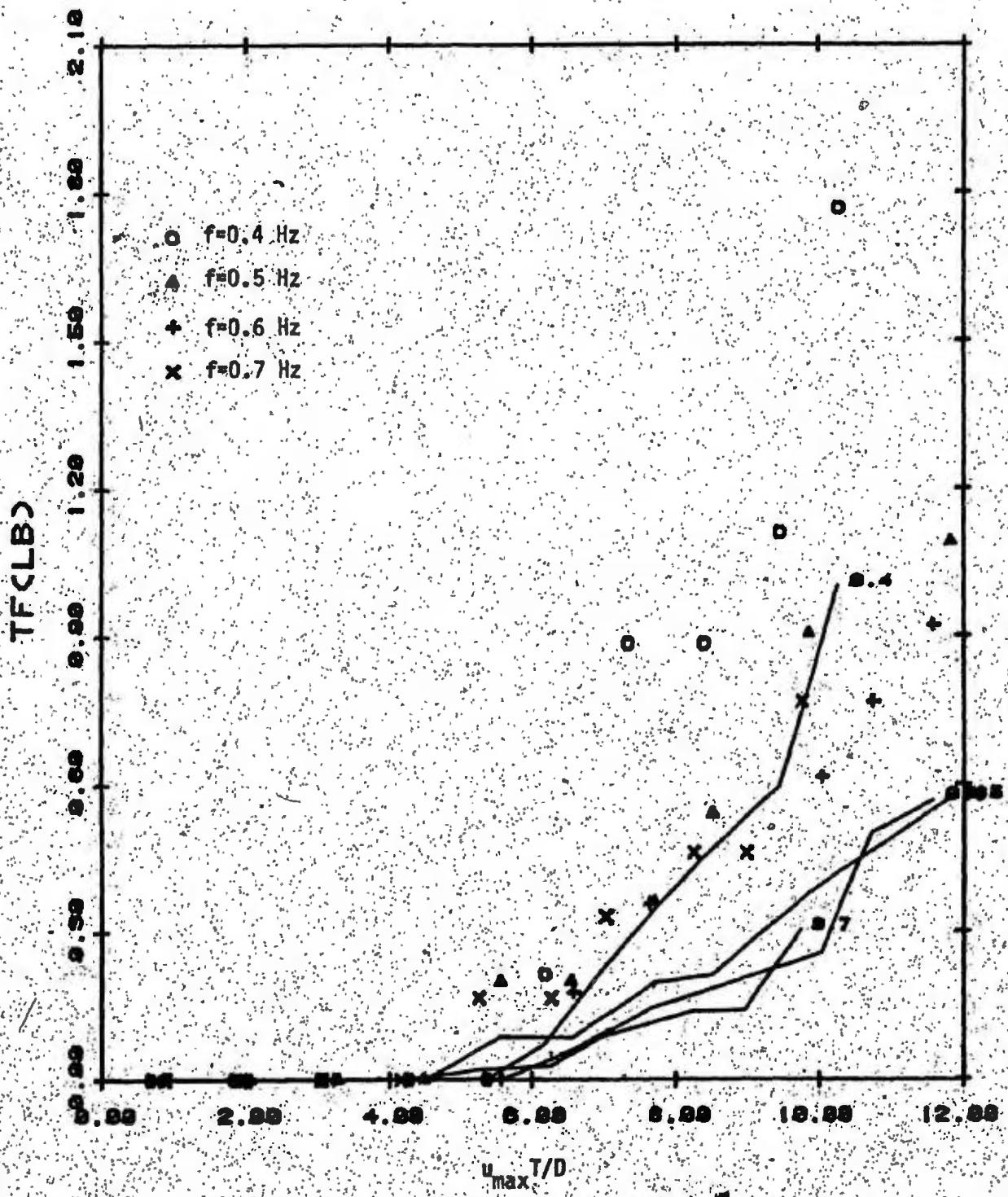


Fig. 29: Transverse Force vs. Keulegan-Carpenter Number

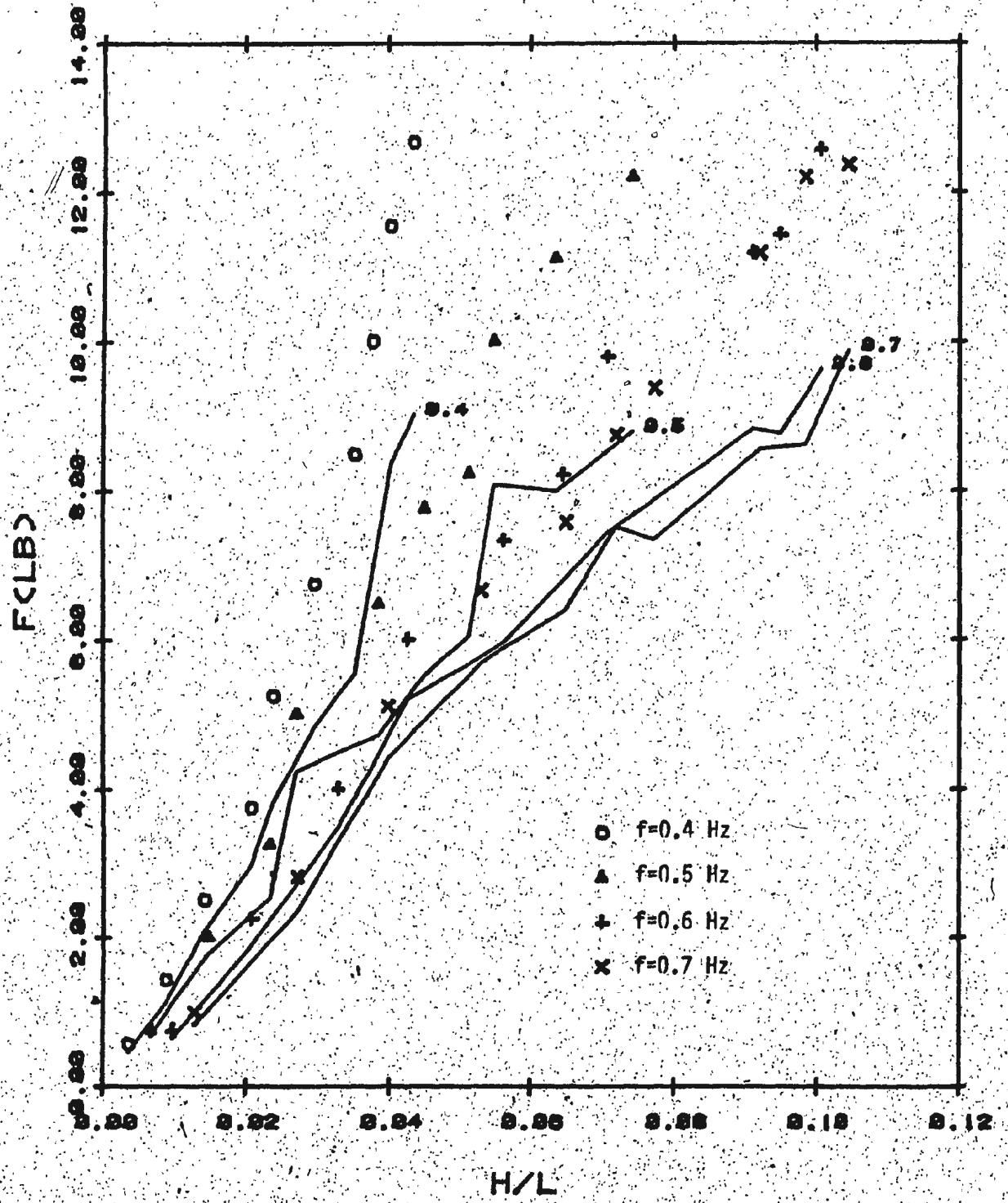


Fig. 30: Longitudinal Force vs. Wave Steepness

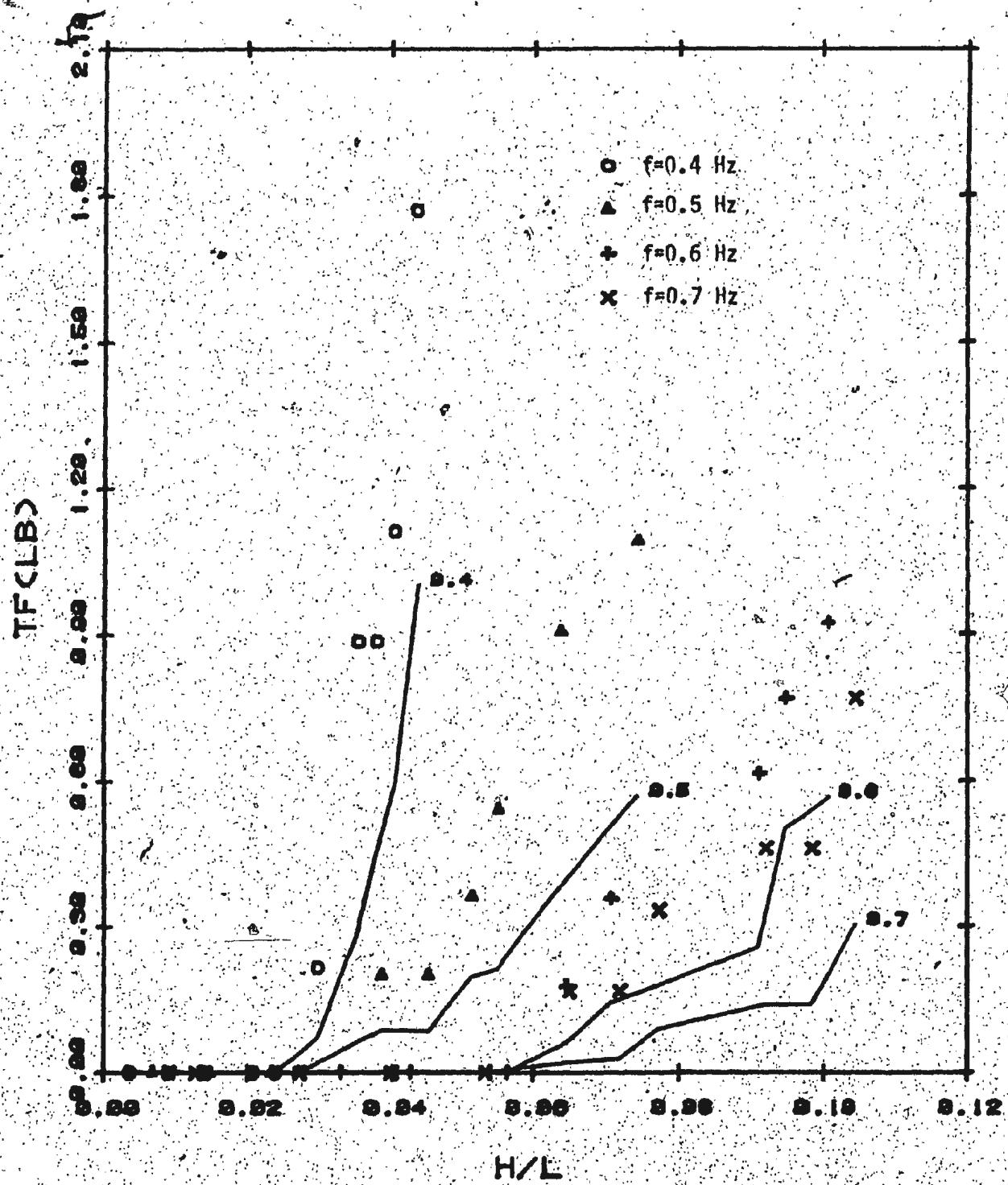


Fig. 31: Transverse Force vs. Wave Steepness

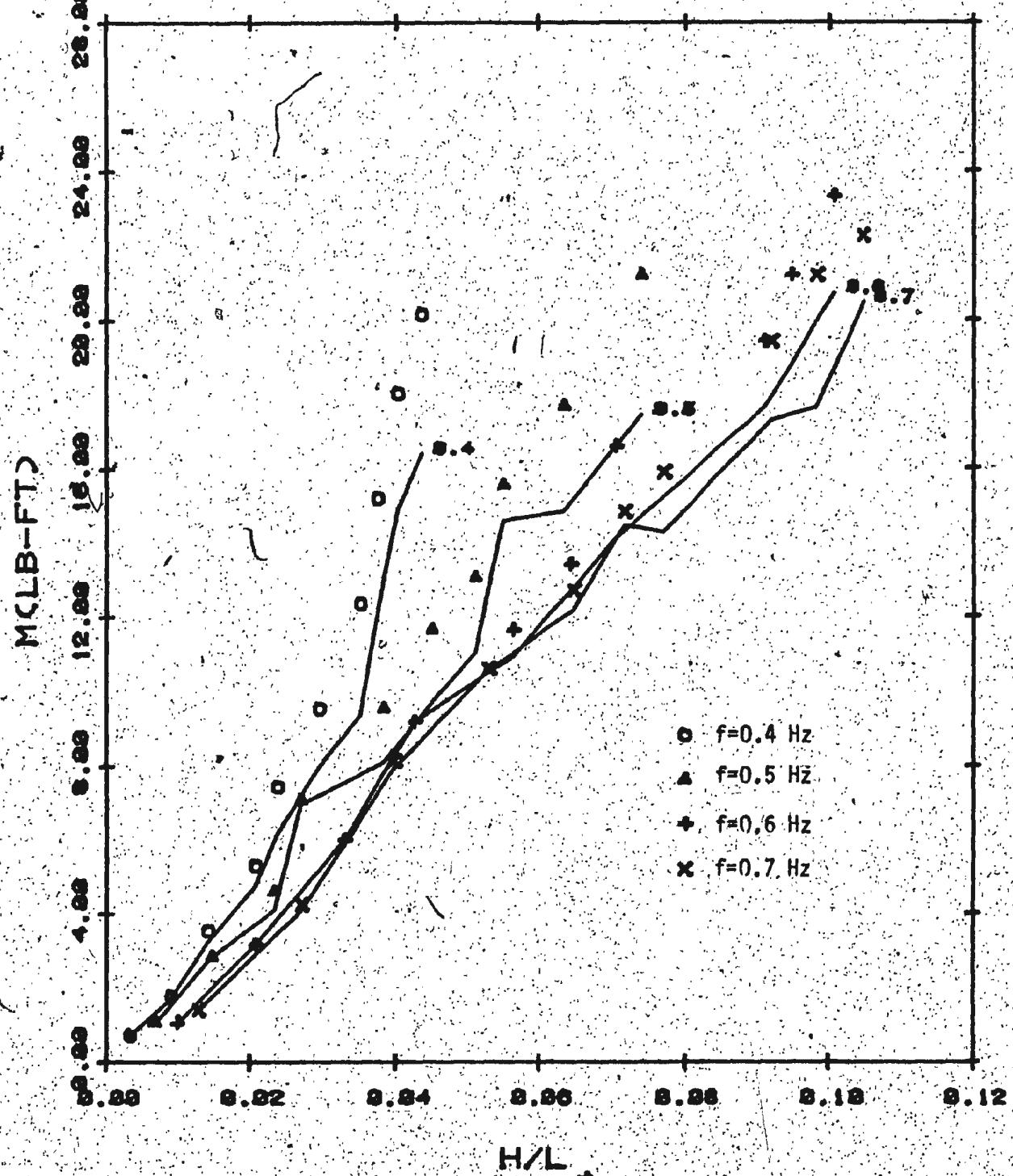


Fig. 32: Longitudinal Moment vs. Wave Steepness

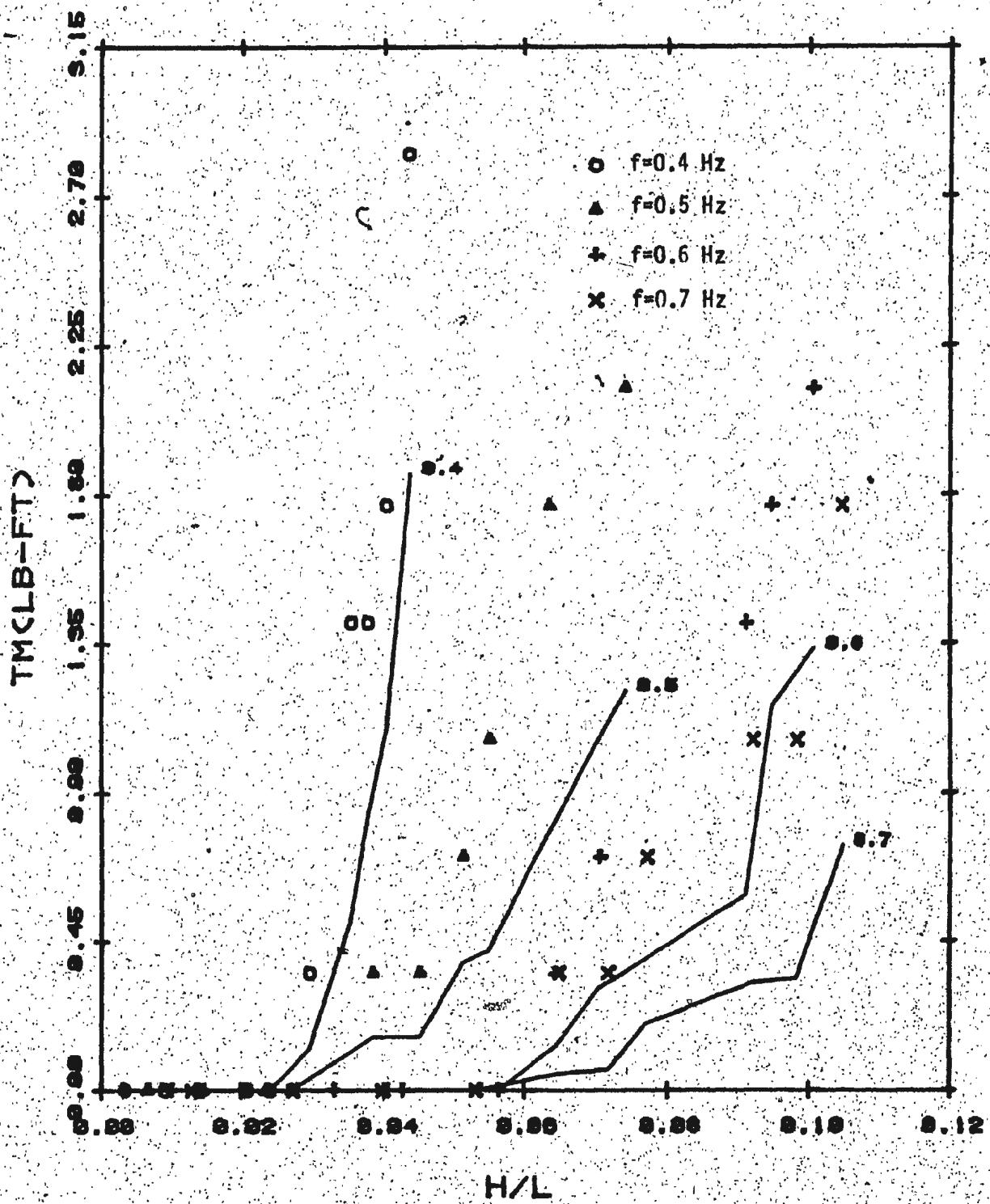


Fig. 33: Transverse Moment vs. Wave Steepness

APPENDIX A

**Computer Program for the Analysis of Piston-Type Wave Generator  
Performance**

```

$COMPILE
C THIS PROGRAM IS BASED ON BIESEL'S THEORY FOR PISTON-TYPE
C WAVE GENERATOR
C H:WAVE HEIGHT(FT)
C L:WAVELENGTH(FT)
C D:WATER DEPTH(FT)
C E:STROKE(INGH)
1 DIMENSION L(120,120),H(120,120),XC(120,120)
2 REAL L
3 D=3.0
4 G=6.2832*D
5 PRINT 12
6 DO 11 I1=5,100,5
7 I=I1/5
8 AI=I1
9 EE=AI/120.
10 E=EE*124
11 OO=2 JJ=10,600,10
12 AJ=JJ
13 J=JJ/10
14 L(I,J)=AJ/10
15 XC(I,J)=D/L(I,J)
16 X=G/L(I,J)*2
17 IF (X.GT.174.672) GO TO 3
18 EX1=EXP(X)
19 EX2=EXP(-1.*X)
20 EX3=EXP(X/2.)
21 H(I,J)=(EX1+EX2-2.)/(EX1-EX2+4.*G/L(I,J))*4.*EE
22 YP=H(I,J)/L(I,J)
23 IF (YP.GT.0.1) GO TO 4
24 GO TO 2
25 H(I,J)=(EX3-2./EX3)/(EX3+4.*G/(L(I,J)*EX3))*4.*EE
26 GO TO 2
27 H(I,J)=0.
28 2 CONTINUE
29 PRINT 10,E
30 10 FORMAT(//,1X,2HE=,F10.4,/)
31 PRINT 13
32 13 FORMAT(3X,1HL,6X,1HH,9(5X,1HE,6X,1HH),//)
33 PRINT 11,(L(I,J),H(I,J),J=1,60)
34 1 CONTINUE
35 11 FORMAT(1D(F5.1,F8.2))
36 PRINT 12
37 12 FORMAT(1H1)
38 STOP
39 END

```

\*EXECUTE

APPENDIX B

Computer Program for the Analysis of Wave Forces and Overturning Moments  
on a Vertical Circular Cylinder Based on Stoke's First Order Wave Theory

\$COMPILE  
 C THIS PROGRAM IS BASED ON THE STOKES' FIRST ORDER WAVE THEORY FOR THE WAVE  
 C FORCES AND MOMENTS ON THE CIRCULAR CYLINDRICAL PILE  
 F:WAVE FREQUENCY(HZ)  
 H:WAVEHEIGHT(FT)  
 L:WAVELLENGTH(FT)  
 D:WATER DEPTH(FT)  
 PD:PILE DIAMETER(FT)  
 REC:RELATIVE DEPTH  
 WS:WAVE STEEPNESS  
 AT:THE RATIO OF PILE DIAMETER TO WAVELENGTH  
 UMAX:THE MAXIMUM HORIZONTAL COMPONENT OF PARTICLE VELOCITY(FT/SEC)  
 UTD:THE KEULEGAN AND CARPENTER PARAMETER  
 RD:REYNOLDS NUMBER  
 NU:KINEMATIC VISCOSITY(FT\*\*2/SEC)  
 CD:DRAG COEFFICIENT  
 CM:MASS COEFFICIENT  
 CL:LIFT COEFFICIENT  
 IN:MOIMENT OF INERTIA(INCH\*\*4)  
 E:YOUNG'S MODULUS(PSI)  
 RO:MASS DENSITY OF WATER(SLUG/FT\*\*3,LB-SEC\*\*2/FT\*\*4)  
 S1:ANY POINT FROM THE BOTTOM OF OCEAN(FT)  
 MF:MEASURING HORIZONTAL FORCE IN THE LONGITUDINAL DIRECTION(LB)  
 MTF:MEASURING HORIZONTAL FORCE IN THE TRANSVERSE DIRECTION(LB)  
 FH:THEORETICAL HORIZONTAL FORCE IN THE LONGITUDINAL DIRECTION(LB)  
 TF:THEORETICAL HORIZONTAL FORCE IN THE TRANSVERSE DIRECTION(LB)  
 EBS:LONGITUDINAL STRAIN(E-06)  
 EBSD:LONGITUDINAL STRAIN(E-06) FOR DRAG COEFFICIENT  
 EBSM:LONGITUDINAL STRAIN(E-06) FOR MASS COEFFICIENT  
 EBT:TRANSVERSE STRAIN(E-06)  
 EBTL:TRANSVERSE STRAIN(E-06) FOR LIFT COEFFICIENT  
 MM:MEASURABLE MOMENT AT POINT S1 IN THE LONGITUDINAL DIRECTION(LB-FT)  
 M1:MEASURABLE MOMENT AT POINT S1 FOR THE DRAG COEFFICIENT  
 M2:MEASURABLE MOMENT AT POINT S1 FOR MASS COEFFICIENT  
 MR:MEASURABLE MOMENT AT POINT S1 IN THE TRANSVERSE DIRECTION(LB-FT)  
 M3:MEASURABLE MOMENT AT POINT S1 FOR THE LIFT COEFFICIENT  
 MT:THEORETICAL MOMENT AT POINT S1 IN THE LONGITUDINAL DIRECTION(LB-FT)  
 TM:THEORETICAL MOMENT AT POINT S1 IN THE TRANSVERSE DIRECTION(LB-FT)  
 BHM:THE PHASE ANGLE OF THE MAXIMUM TOTAL MOMENT WILL BE OCCURRED(DEGREES)  
 BHF:THE PHASE ANGLE OF THE MAXIMUM TOTAL HORIZONTAL FORCE WILL BE  
 OCCURRED(DEGREES)  
 L1:THE LOCATION OF THE APPLICATION OF THE RESULTANT HORIZONTAL FORCE IN  
 THE LONGITUDINAL DIRECTION(FT)  
 L2:THE LOCATION OF THE APPLICATION OF THE RESULTANT HORIZONTAL FORCE IN  
 THE TRANSVERSE DIRECTION(FT)  
 1 DIMENSION UTD(40),RD(40),F(40),H(40),L(40),MM(40),MT(40),UMAX(40),  
 6FH(40),BHM(40),BHF(40),CD(40),CM(40),L1(40),EBS(40),IK(40),  
 5RED(40),WS(40),AT(40),AF(40),EBT(40),MR(40),CL(40),TF(40),TM(40),  
 2L2(40),MF(40),MTP(40),EBSD(40),EBSM(40),EBTL(40),M1(40),M2(40),  
 3M3(40)  
 2 REAL L,NU,MM,MT,IN,K1,K2,K3,K4,L1,MR,L2,MF,MTP,M1,M2,M3  
 3 C=3.  
 4 FD=0.547  
 5 NU=0.00001  
 6 RO=2.  
 7 PI=3.1416  
 8 IN=27.41  
 9 E=507300.  
 10 C=3.28125  
 11 S1=0.23

```

12      CO 15 I=1,40
13      15 FREAD(5,7) MF(I),MTF(I)
14      DO 10 I=1,40
15      IK(I)=I
16      READ(5,1) F(I),H(I),L(I),EBS(I),EBT(I)
17      READ(5,50) EBSD(I),EBSM(I),EBTL(I)
18      S2=H(I)/2.+D
19      X1=2.*PI*S2/L(I)
20      X2=2.*PI*D/L(I)
21      PDC(I)=D/L(I)
22      WS(I)=H(I)/L(I)
23      AT(I)=PD/L(I)
24      AF(I)=D*F(I)**2
25      COSHI=(EXP(X1)+EXP(-X1))/2.
26      SINH1=(EXP(X2)-EXP(-X2))/2.
27      UMAX(I)=PI*H(I)*F(I)*COSHI/SINH1
28      UTD(I)=UMAX(I)/(F(I)*PD)
29      RD(I)=UMAX(I)*PD/NU
30      R1=2.*PI*S1/L(I)
31      P1_2=2.*R1
32      P2_2=2.*X1
33      SINH2=(EXP(X1)-EXP(-X1))/2.
34      SINH3=(EXP(R22)-EXP(-R22))/2.
35      SINH4=(EXP(R1)-EXP(-R1))/2.
36      SINH5=(EXP(R12)-EXP(-R12))/2.
37      COSH2=(EXP(X1)+EXP(-X1))/2.
38      COSH3=(EXP(R22)+EXP(-R22))/2.
39      COSH4=(EXP(R1)+EXP(-R1))/2.
40      COSH5=(EXP(R12)+EXP(-R12))/2.
41      K1=(R22-R12+SINH3-SINH5)/(16.*SINH1**2)
42      K2=(SINH2-SINH4)/SINH1
43      K3=(1./(64.*SINH1**2))*(0.5*(P22)**2-0.5*(R12)**2+R22*SINH3-R12*
7SI NH5-COSH3+COSH5)
44      K4=(1./(2.*SINH1))*(X1*SINH2-R1*SINH4-COSH2+COSH4)
45      MM(I)=E*EBS(I)*IN/(12.*C)
46      M1(I)=E*EBSD(I)*IN/(12.*C)
47      M2(I)=E*EBSM(I)*IN/(12.*C)
48      CD(I)=M1(I)/(RD*PD*(F(I)*H(I)*L(I))**2*(K3-R1*K1/2.))
49      CM(I)=(4.*M2(I))/(PI*H(I)*RD*(F(I)*L(I)*PD)**2*(K4-R1*K2/2.))
50      PMC=ABS(PI*PD*CM(I)*(8.-P1*K2/2.)/(8.*H(I)*CD(I)*(K3-R1*K1/2.)))
51      IF(PMC.GT.1.) GO TO 20
52      GO TO 21
53      20 PMC=1.
54      21 BHMR=ARSIN(PMC)
55      BHM(I)=180./PI*BHMR
56      PMD=ABS(PI*PD*CM(I)*K2/(8.*H(I)*CD(I)*K1))
57      IF(PMD.GT.1.) GO TO 22
58      GO TO 23
59      22 PMD=1.
60      23 BHFR=ARSIN(PMD)
61      BHF(I)=180./PI*BHFR
62      FH(I)=PI*RD*PD*(F(I)*H(I))**2*L(I)*(PI*PD/(4.*H(I))*CM(I)*K2*SIN(
2*BHFR)+CD(I)*K1*ABS(COS(BHFR))*COS(BHFR))
63      MT(I)=RD*PD*(F(I)*H(I)*L(I))**2*(PI*PD/(4.*H(I))*CM(I)*K4*SIN(BHMF
3)+CD(I)*K3*ABS(COS(BHMR))*COS(BHMR)-R1*(PI*PD/(8.*H(I))*CM(I)*
4K2*SIN(BHMR)+0.5*CD(I)*K1*ABS(COS(BHMR))*COS(BHMR)))
64      L1(I)=MT(I)/FH(I)
65      NR(I)=E*EBT(I)*IN/(12.*C)
66      N3(I)=E*EBTL(I)*IN/(12.*C)
67      CL(I)=M3(I)/(RD*PD*(F(I)*H(I)*L(I))**2*(K3-R1*K1/2.))

```

```

68      TF(I)=PI*RD*PD*(F(I)*H(I))**2*L(I)*CL(I)*K1
69      TM(I)=RD*PD*(F(I)*H(I)*L(I))**2*CL(I)*(K3-F1*K1/2.)
70      IF(TF(I).EQ.0.) GO TO 12
71      L2(I)=TM(I)/TF(I)
72      GO TO 10
73      12 L2(I)=0.
74      10 CONTINUE
75      WRITE(6,2)
76      DO 5 I=1,40
77      5 WRITE(6,3)IK(I),F(I),H(I),L(I),RD(I),WS(I),AT(I),AF(I),UMAX(I),
    BUTD(I),RD(I),MM(I),CD(I),CM(I)
78      WRITE(6,9)
79      DO 6 I=1,40
80      6 WRITE(6,11)IK(I),BHM(I),BHF(I),MT(I),FH(I),MF(I),L1(I),MR(I),
    9CL(I),TF(I),MTF(I),TM(I),L2(I)
81      WRITE(6,68)
82      1 FORMAT(F5.1,2F10.3,2E15.2)
83      50 FORMAT(3E15.2)
84      2 FORMAT(1H1,6X,1HF,5X,1HH,6X,1HL,8X,3HD/L,5X,3HH/L,5X,4HPD/L,4X,6HD
    1*F**2,1X,4HUMAX,8X,3HUTD,9X,2HRD,11X,2HMM,5X,2HCD,7X,2HCM,//)
85      3 FORMAT(1S,F5.1,6F8.3,3E12.4,3F8.3)
86      7 FORMAT(2F10.3)
87      9 FORMAT(1H1,8X,3HBHM,7X,3HBHF,7X,2HMT,8X,2HFH,8X,2HMF,9X,2HL1,8X,
    12HMR,8X,2HCL,8X,2HTF,8X,3HMTF,7X,2HTM,8X,2HL2,//)
88      11 FORMAT(1S,12F10.3)
89      68 FORMAT(1H1)
90      STOP
91      END

```

SEXECUTE

